

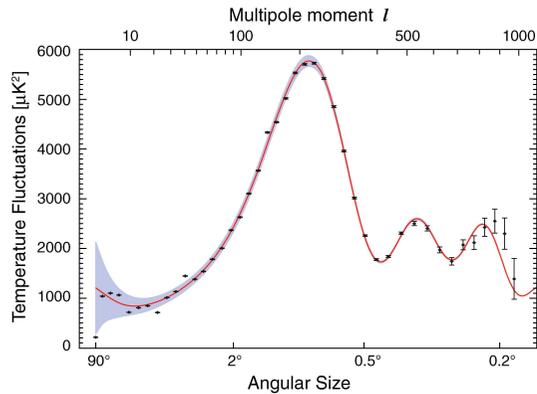


Combining direct detection with colliders

Based on the paper (in preparation) by G. Bertone, D. G. Cerdeno,
M.F., R. Ruiz de Austri and R. Trotta

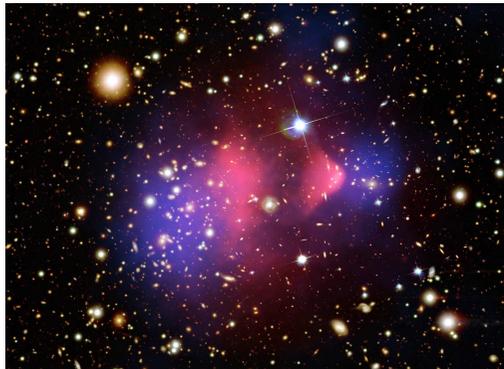
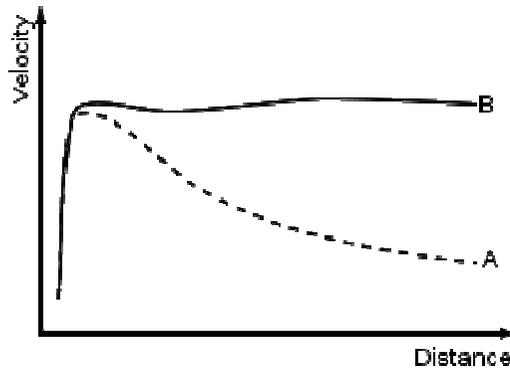
Mattia Fornasa
University of Padova – INFN Padova

Dark Matter and experimental techniques

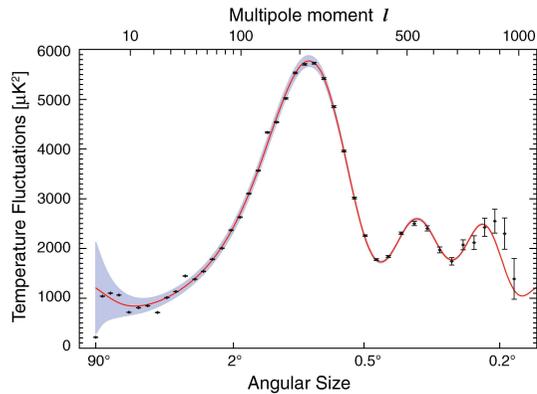


DM evidences:

- angular spectrum of anisotropies in the CMB
- rotation curves of galaxies
- weak lensing mass reconstruction for interacting clusters of galaxies



Dark Matter and experimental techniques

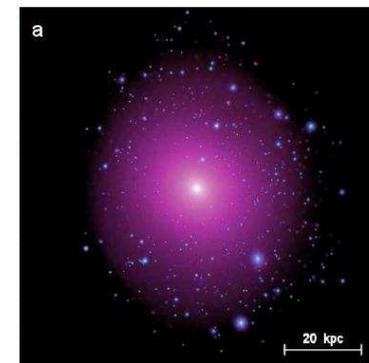
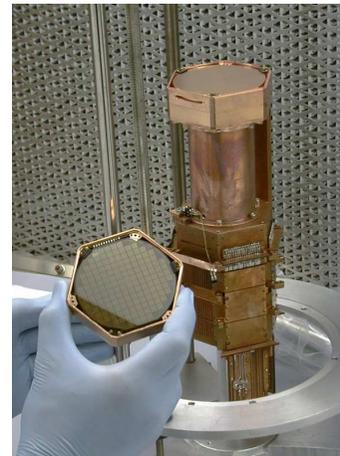
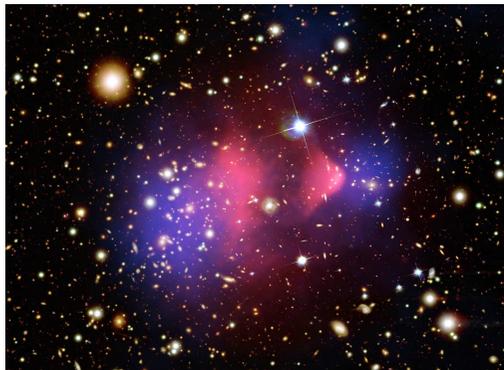
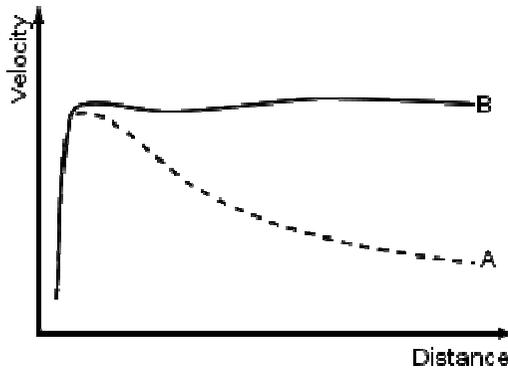


DM evidences:

- angular spectrum of anisotropies in the CMB
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Additional evidences and future detections are expected from:

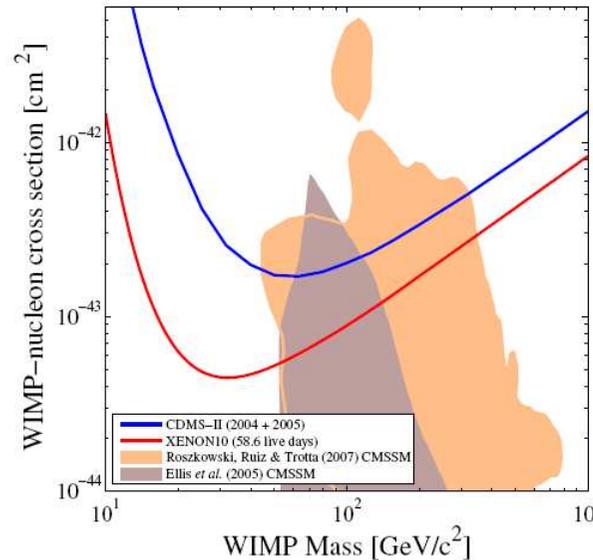
- **colliders**
- **direct detection**
- **indirect detection**



DM direct detection

Based on the possibility of detecting the **recoil energy** deposited by a DM particle to the nuclei of the detector when the particle passes through the detector itself.

CDMS-II detection of 2 events with an estimated background of 0.9



From XENON10 Collaboration, Phys. Rev. Letters, 100, 021303 (2008)

Predicted **event number**:

$$\lambda = \epsilon \int_{E_{th}}^{E_{max}} \frac{dR}{dE}(E) F^2(E) dE = \epsilon \int_{E_{th}}^{E_{max}} c_1 R_0 e^{-E/(c_2 E_0)} F^2(E) dE \quad (1)$$

$$R_0 = \frac{\sigma_{\chi,p}^{SI} \rho_{\chi} A^2 c^2 (m_{\chi} + m_p)^2}{\sqrt{\pi} m_{\chi}^3 m_p^2 v_0} \quad E_0 = \frac{2m_{\chi}^2 v_0^2 A m_p}{(m_{\chi} + A m_p)^2 c^2} \quad (2)$$

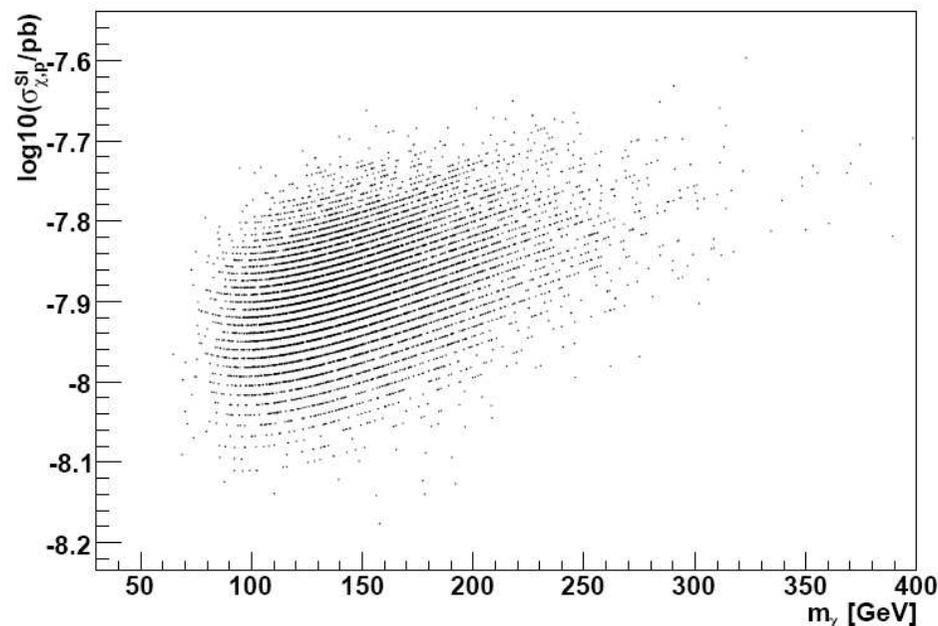
Reconstructing DM properties: direct detection

Reconstruction from direct detection has been studied by Anne Green (hep-ph/0703217 and 0805.1704):

- choose a **benchmark model**
- predict the response of a direct detection experiment $(\lambda, \{\varepsilon_i\})$
- simulate a large number of experiments $(N, \{E_i\})$
- associate a **likelihood** function to each of them

$$\mathcal{L} = \frac{e^{-N\lambda} \lambda^N}{N!} \prod_{i=1}^N f(E_i) \quad (3)$$

- **Maximum-likelihood estimators** for m_χ and $\sigma_{\chi,p}^{\text{SI}}$ can be obtained imposing the constraints $\partial L/\partial m=0$ and $\partial L/\partial \sigma=0$
- Constant and exponential **background** are studied



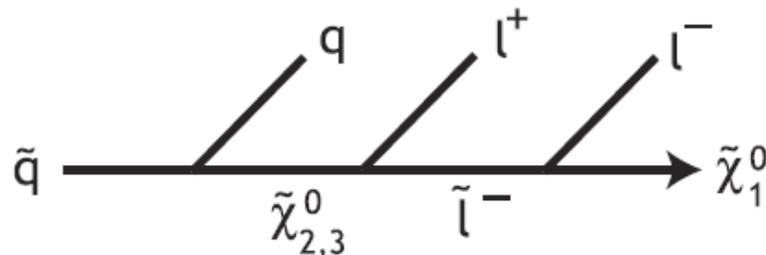
Colliders phenomenology - I

LHC (finally restarted!) will access an unexplored range of energies, with the possibility of detecting new, heavy particles beyond the Standard Model.

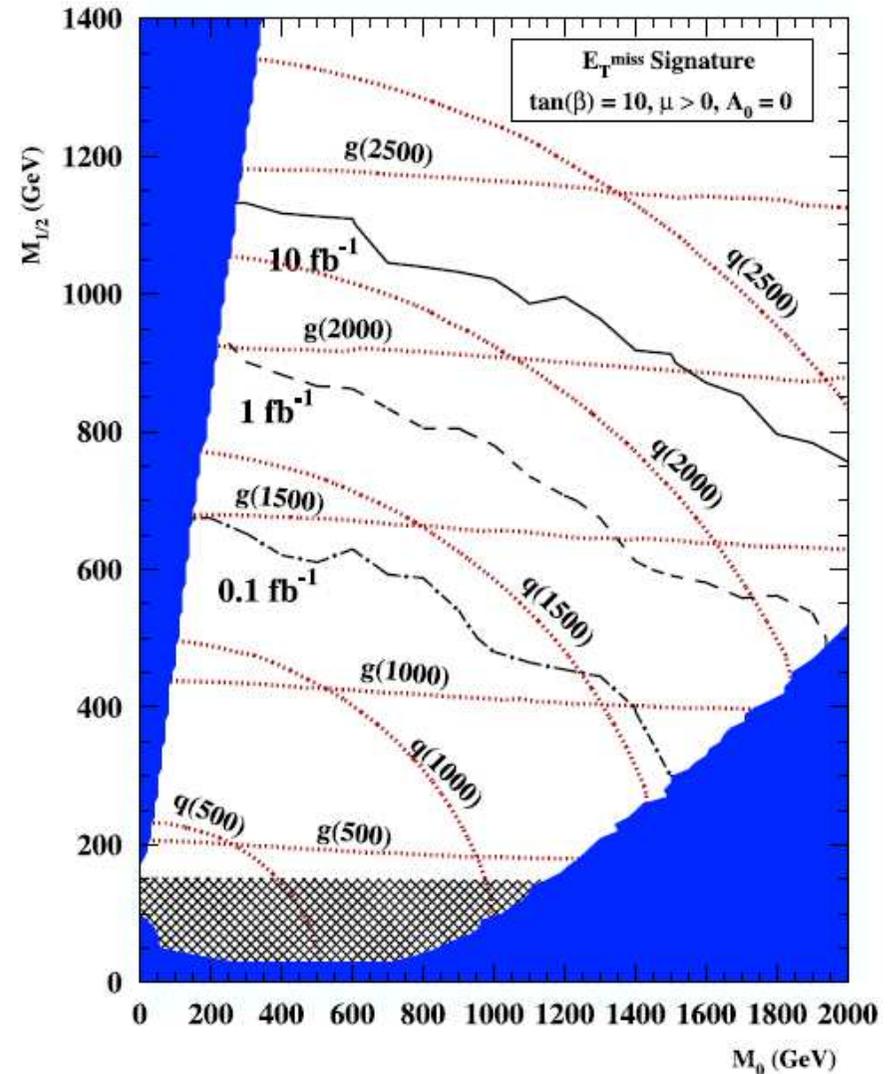
WIMP DM candidate leaves the detector appearing as **unbalanced, missing energy**.

Possibility of measuring other non-SM particles:

- they decay into the WIMP
- theoretical model can be constrained more efficiently



From Tovey, Eur. Phys. J.,
direct C4, N4

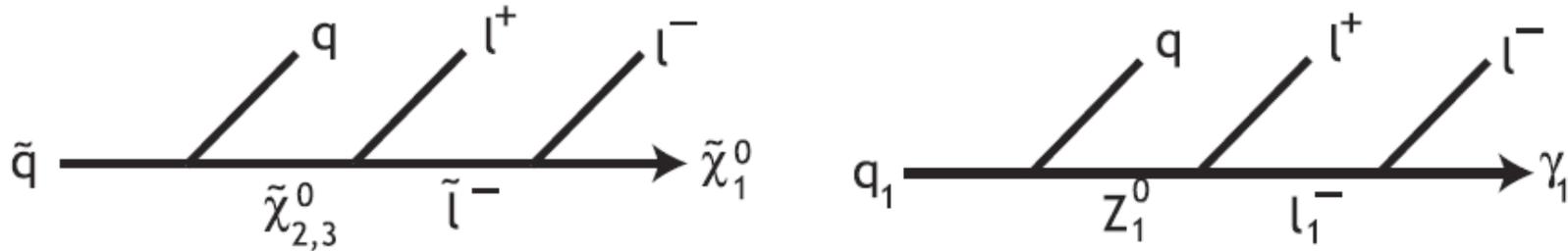


From Baltz et al., Phys. Rev. D74,
103521 (2006)

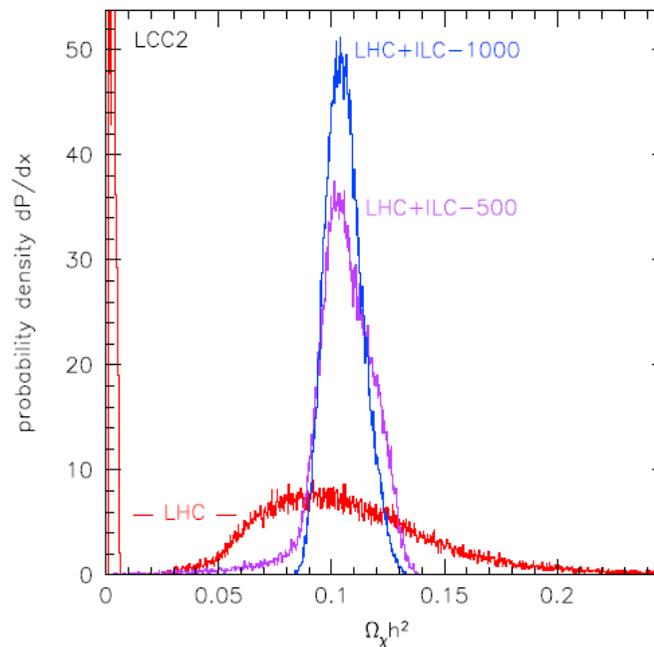
Colliders phenomenology - II

Difficult and challenging task:

- **degeneracies** among theoretical models



- **degeneracies** among parameters



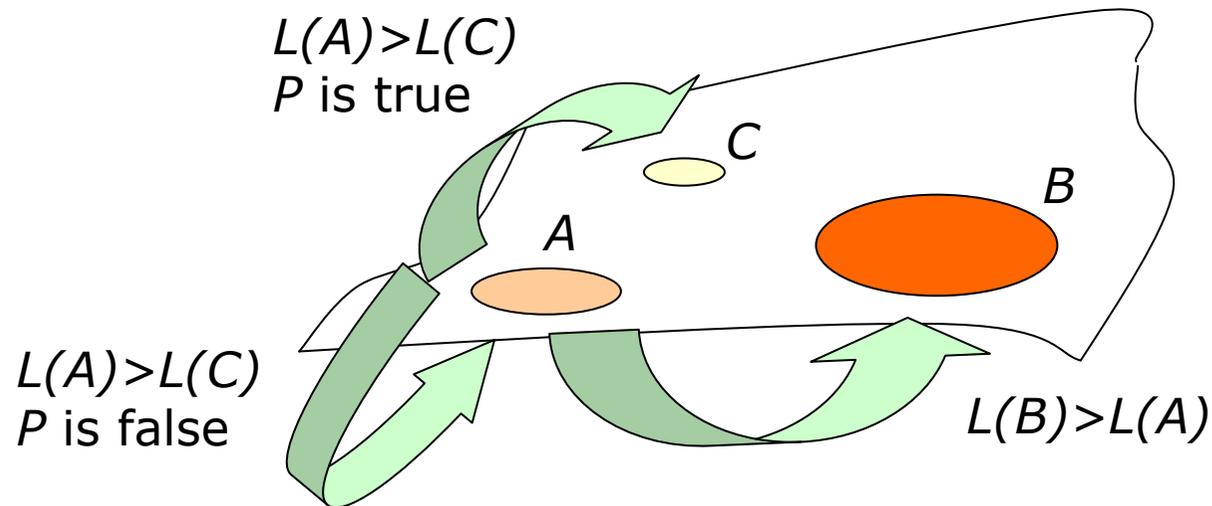
From Baltz et al.,
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Monte Carlo Markov Chains

First input is a set of experimental measurements $\{d_i, \sigma_i\}$.

SUSY parameters space is scanned with the use of **MCMCs** (based on the **Bayesian theorem**) and a likelihood is associated to each point:

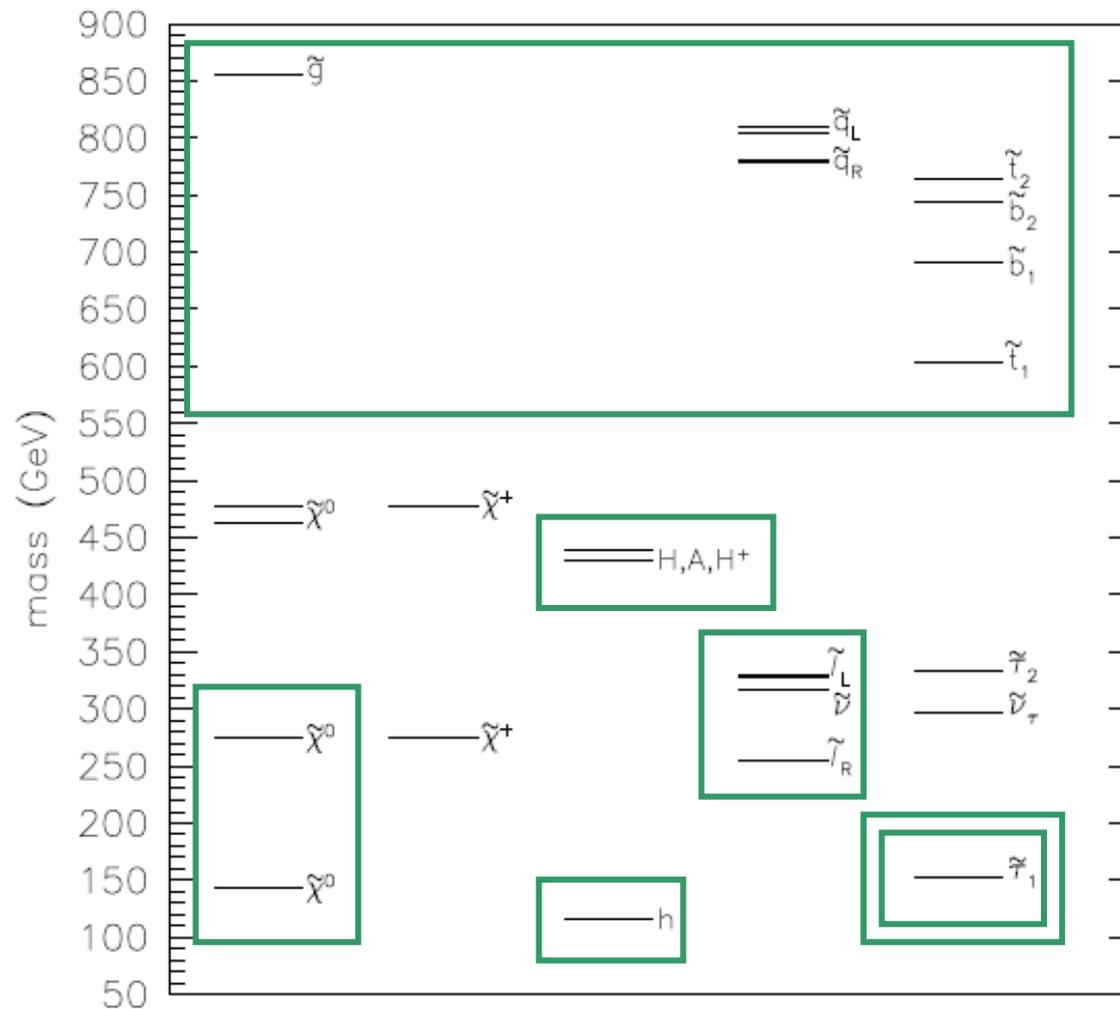
$$\mathcal{L} = \prod_{i=1}^M \exp \left(-\frac{(d_i - d_i^{(j)})^2}{2\sigma_i^2} \right) \quad (4)$$



Posterior probability distribution function (**pdf**) of physical observables $(m_{\chi}, \sigma_{\chi, p}^{\text{SI}}, \Omega_{\chi} h^2)$ is obtained by counting the **multiplicity** within the chains.

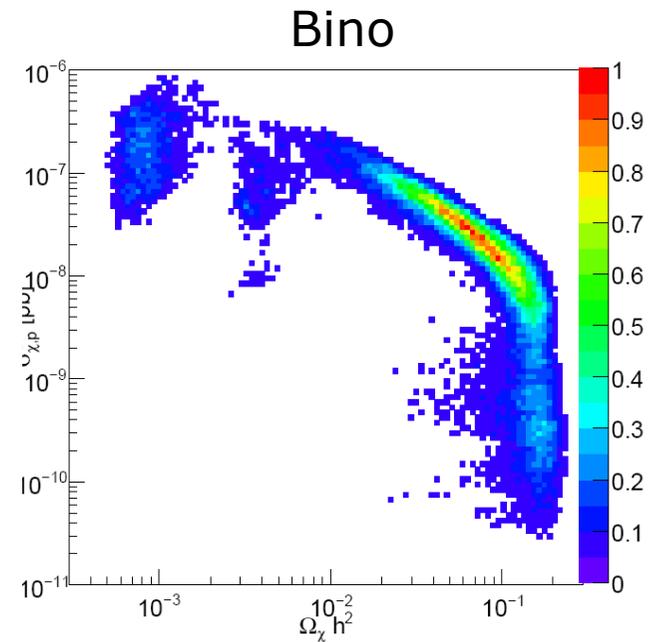
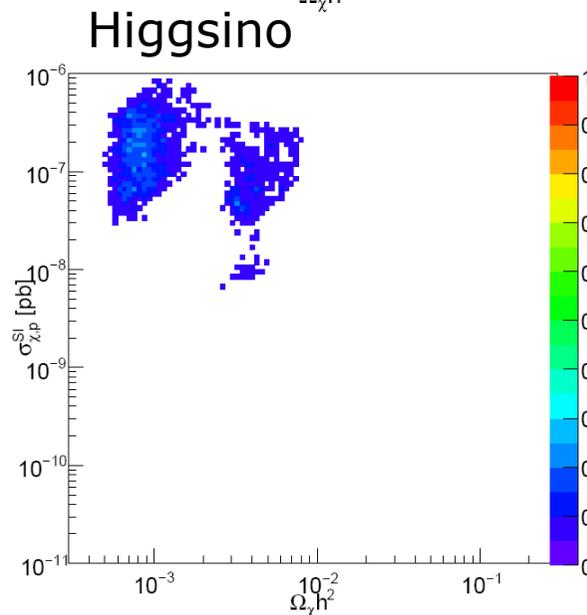
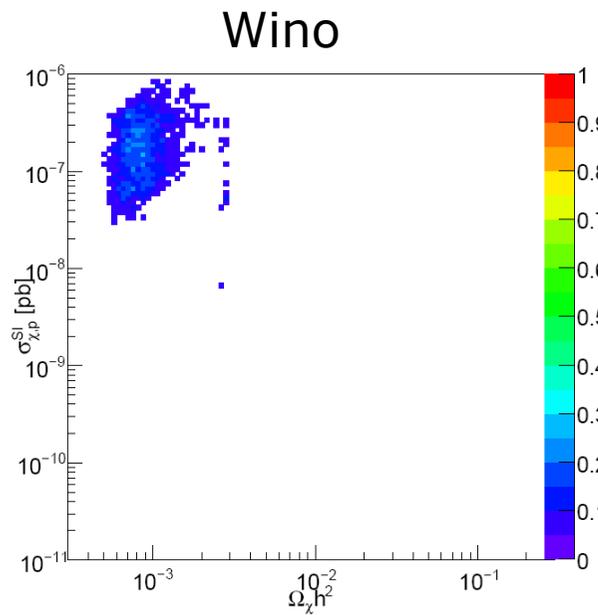
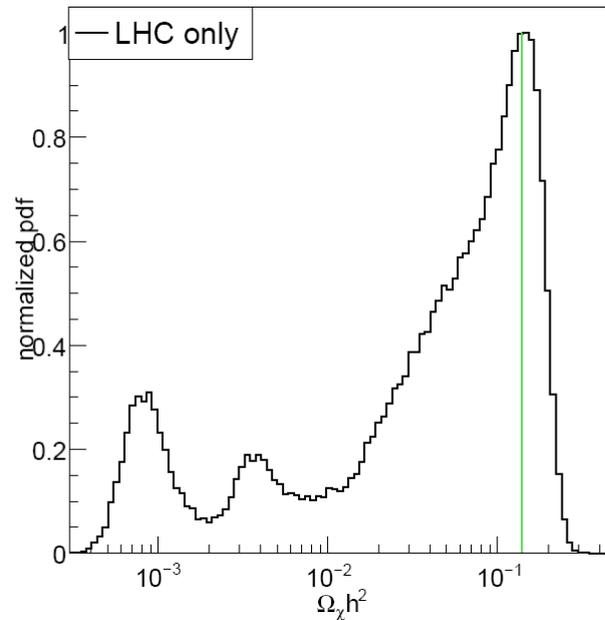
Reconstructing DM properties: colliders

- SUSY parameters space is 24 dimensional
- our benchmark model is in the coannihilation region



From Baltz et al.,
Phys. Rev. D 74 (2006)
103521

Reconstructing DM properties: $\Omega_\chi h^2$



Combining colliders with direct detection

MCMCs can be **sampled** in order to account for informations from direct detection, i.e. the multiplicity of each point is changed by a factor:

$$m_i \longrightarrow m_i \exp\left(-\frac{(\lambda - n^{(i)})^2}{2n^{(i)}}\right) \prod_{\text{bins}, j=1}^{10} \exp\left(-\frac{(n_j - n_j^{(i)})^2}{n_j^{(i)}}\right) \quad (5)$$

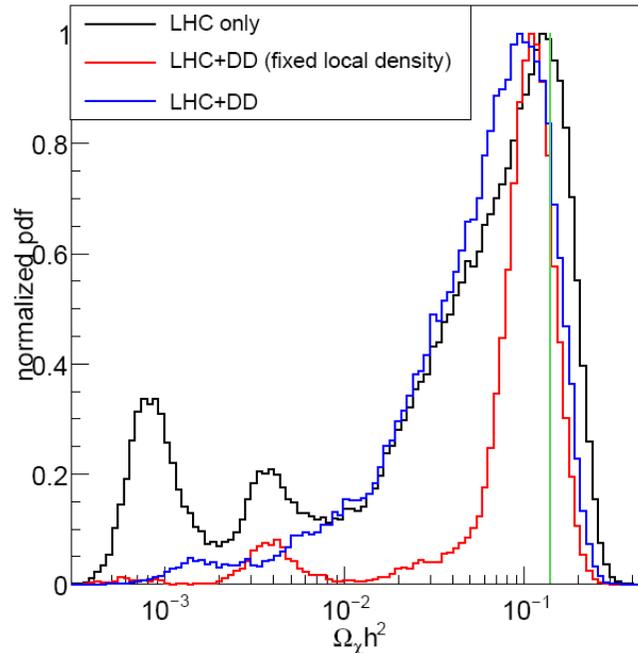
Relic density $\Omega_\chi h^2$:

- assuming a signal at LHC
- assuming that **the same particle** leaves a signal in a direct detection experiment

From the reconstruction of $\Omega_\chi h^2$ (breaking of degeneracies) it is possible to **identify** that particle as the cosmological DM (comparison with WMAP value).

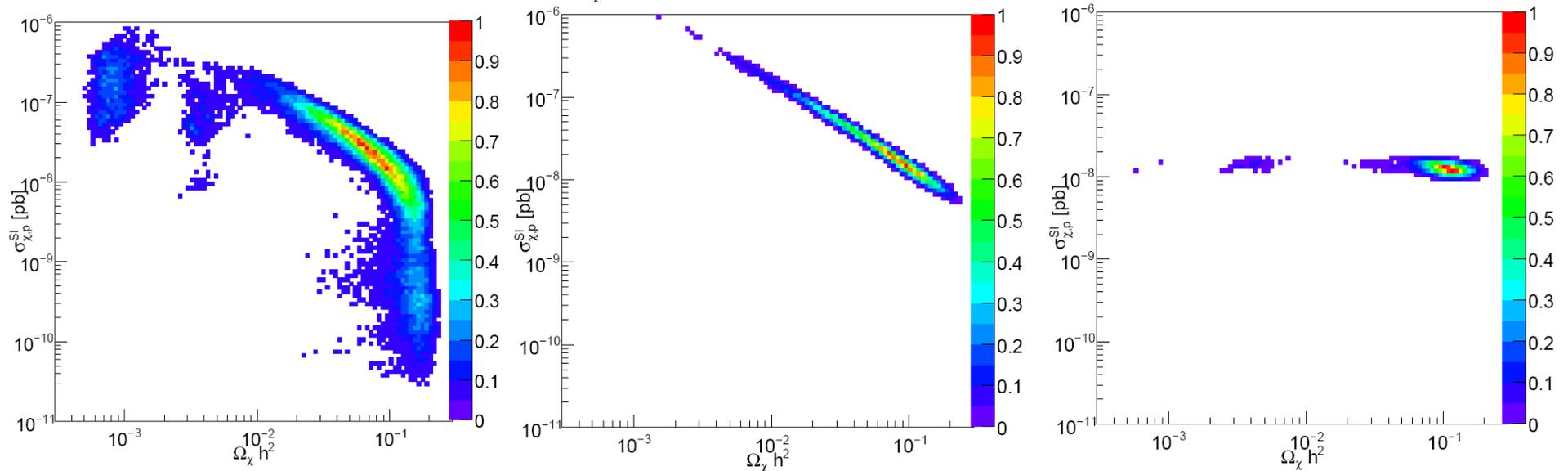
Local density should be rescaled in the case of multi-component DM by a factor $\Omega_\chi^{(i)}/\Omega_\chi^{\text{WMAP}}$.

Results



LHC only: $\Omega_\chi h^2 = 0.147_{0.035}^{0.206}$
(76%)

LHC+DD: $\Omega_\chi h^2 = 0.109_{0.035}^{0.165}$
(68%)



Conclusions

- Direct detection provides a good reconstruction of $\sigma_{\chi, P}^{\text{SI}}$
- LHC can constrain DM observables with the use, e.g., of MCMCs
- Internal degeneracies
- Combination of the two experimental techniques may largely improve the situation
- Breaking the degeneracies for the reconstruction of $\Omega_{\chi} h^2$
- The particle detected at LHC that, at the same time, leaves a signal in a direct detection experiment, can be identified as the DM and
- LHC may be used as a DM experiment

Neutralino nature

m_1 , m_2 and μ are the parameters that determines the nature of neutralino.

Our benchmark model has $m_1 < m_2 < \mu$, but the fact that only the two lightest neutralinos are measured leads to models with other hierarchies that equally fit well the data:

Neutralino nature

