Overview of String Phenomenology

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Outline

- Motivation/Introduction/Generalities
- 🗳 Heterotic models
- D-brane models, F-theory models
- Moduli stabilization and supersymmetry breaking
- Solutions Providence of the second se

String Phenomenology

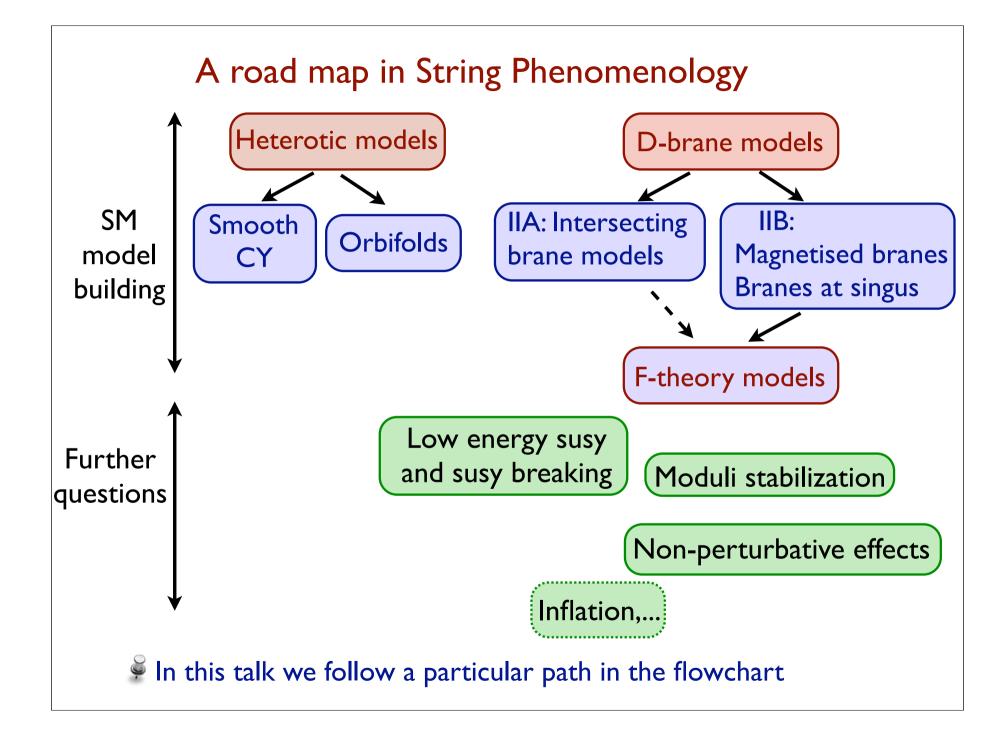
String theory describes gravitational and gauge interactions in a unified framework, consistent at the quantum level

If string theory is realized in Nature, it should be able to describe a very specific gauge sector: Standard Model

- Aim of String Phenomenology:
- Determine classes of constructions with a chance to lead to SM Non abelian gauge interactions, replicated charged fermions, Higgs scalars with appropriate Yukawa couplings, ...

- Within each class, obtain explicit models as close to SM as possible with the hope of learning more about the microscopics of SM in string theory

Old program, yet continuous progress
Moduli stabilization, non-perturbative effects, ...

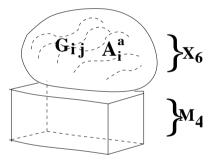


Heterotic string models

[Candelas, Horowitz, Strominger, Witten, '85]

Free IOd heterotic string has as effective theory IOd N=I sugrations coupled to E8xE8 (or SO(32)) gauge multiplets

© Compatification: six extra dimensions parametrize small Calabi-Yau space, on which we also turn on a non-trivial gauge field background

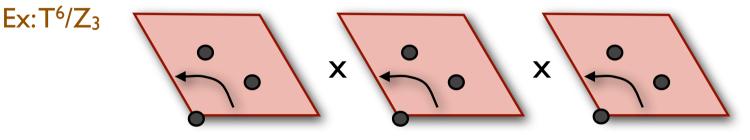


Gauge group is broken by gauge background
Possible to break down to something close to SM gauge group
4d charged chiral fermions from KK reduction of 10 gauginos.
Number of families fixed by topological data

Within this general class, very explicit models close to (MS)SM [e.g. Ovrut et al.]

Heterotic String Orbifolds

A very tractable version is provided by toroidal orbifold models $\stackrel{\circ}{\Rightarrow}$ Compatification space is quotient of T⁶ by discrete symmetry



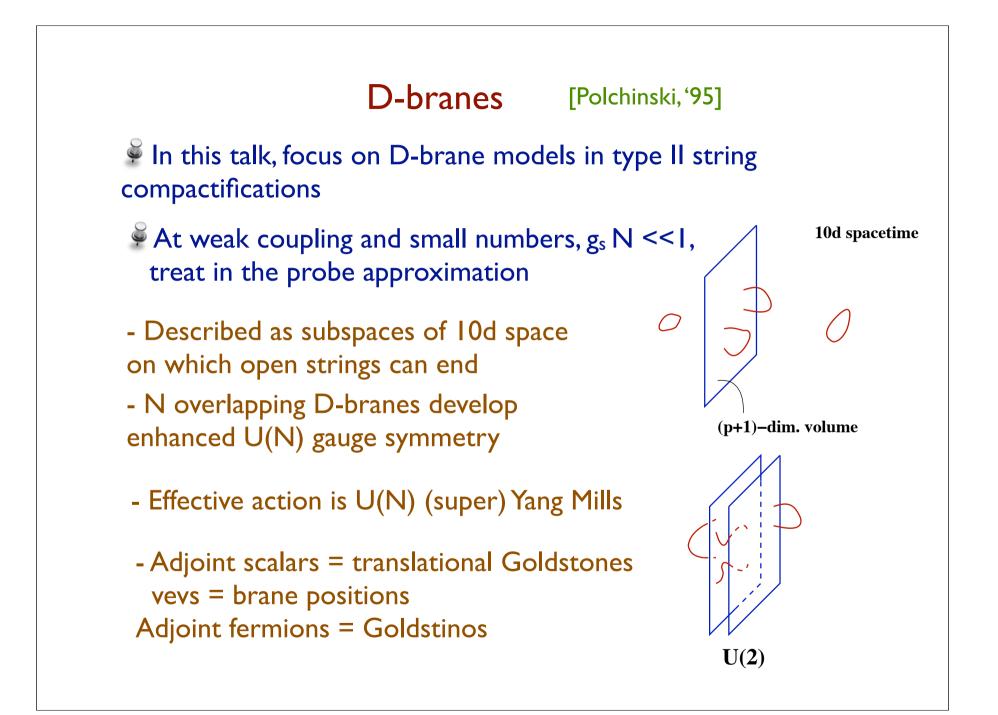
Gauge background is trivial except around orbifold singularities Discrete gauge transformation while moving around fixed point

Models very close to (MS)SM have been explicitly constructed and studied [Talks by Hebecker, Ratz, ...]

- Geometric intuitions very helpful:

Partial gauge breaking at orbifold points/planes

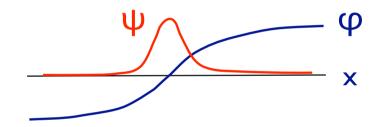
- Need effective field theory arguments to remove exotics, break some gauge symmetries



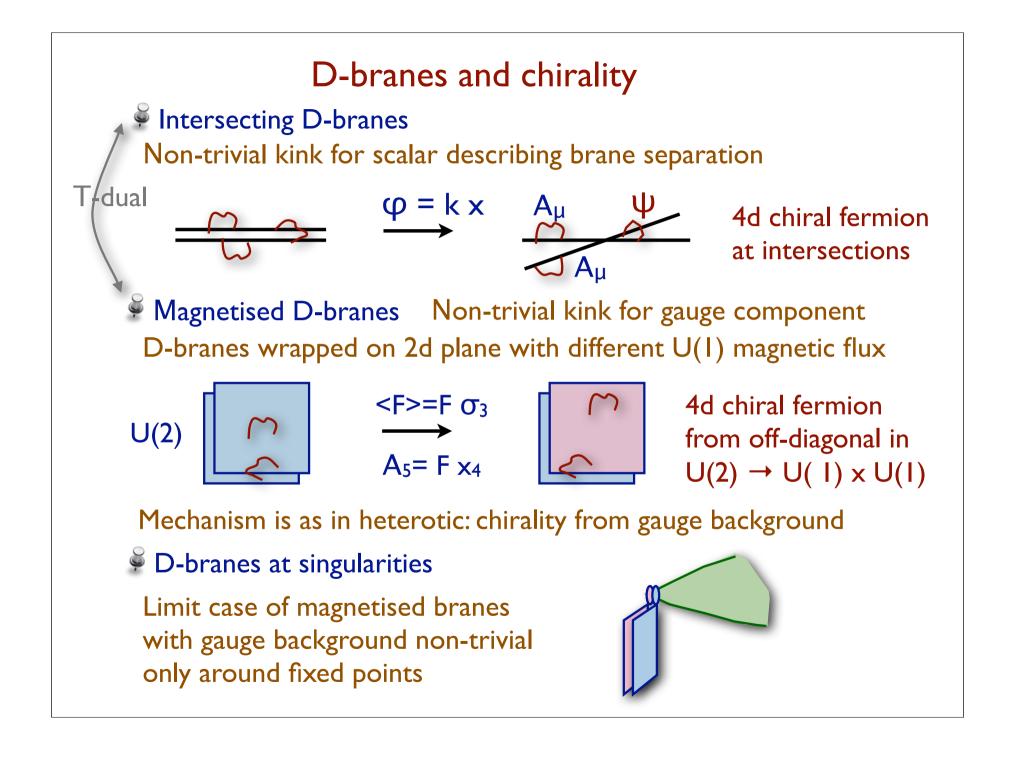
Chirality

Chirality is a crucial criterion discriminating which string/brane configurations can lead to models of particle physics

Mechanism for chirality in string th. can be understood in field th.
5d fermion coupled to scalar kink has 4d chiral fermion zero mode



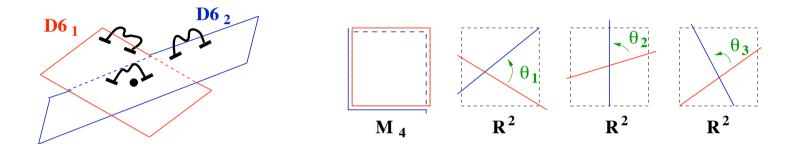
- Domain wall fermions in lattice gauge field theories
- Fermion localization in extra dimensions in BSM phenomenology
- Different string realizations lead to different string constructions
- Different string realizations are dual to each other Not that "different"!



Intersecting D6-branes in type IIA

[Berkooz, Douglas, Leigh,'96]

Consider type IIA string theory with two stacks of D6-branes (hence 7d subspaces) intersecting over a 4d subspace of their volumes



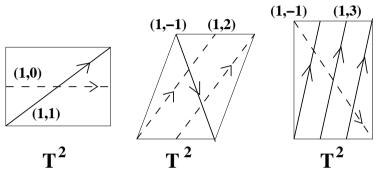
Three sectors of open strings

- D61-D61: U(N1) on 7d plane 1
- D62-D62: U(N2) on 7d plane 2
- D61-D62: 4d chiral fermion in (N_1, \overline{N}_2) on 4d intersection
- Chirality is a consequence of the geometry of the intersection e.g. two D5's intersecting over 4d leads to non-chiral fermions Need intersections in all three complex planes

Intersecting brane worlds

[Blumenhagen, Gorlich, Kors, Lust; Aldazabal, Franco, Ibanez, Rabadan, AU; '00]

 $\stackrel{\scriptstyle \eq}{=}$ To obtain 4d gravity, need to compactify, e.g. on on T⁶=T² xT² xT²



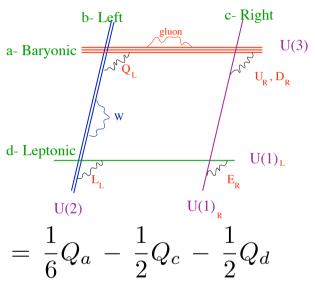
Solutions of D6-branes in sets of N_a D6_a-branes wrapping 3cycles Π_a described as products of 1-cycles $(n_a{}^i, m_a{}^i)$ on each $(T^2)^i$

Gauge group $\prod_{a} U(N_a)$ Chiral fermions $\sum_{a,b} I_{ab} (N_a, \overline{N}_b)$ $I_{ab} = (n_a^1 m_b^1 - n_b^1 m_a^1) \times (n_a^2 m_b^2 - n_b^2 m_a^2) \times (n_a^3 m_b^3 - n_b^3 m_a^3)$ Intersection number = geometric origin of family replication! \checkmark Generalizes to D6-branes on intersecting 3-cycles on general CY

Towards the SM

A simple road to SM (not unique, e.g. see later for GUTs) Introduce four stacks of D6's a,b,c,d $U(3)_a \times USp(2)_b \times U(1)_c \times U(1)_d$ $I_{ab} = 3 \rightarrow Q_L$ $I_{ac} = -3, I_{ac'} = 3 \rightarrow U_R, D_R$ $I_{db} = 3 \rightarrow L$ $I_{dc} = -3, I_{dc'} = -3 \rightarrow E_R, \nu_R$ Spectrum of SM with hypercharge $Y = \frac{1}{6}Q_a - \frac{1}{2}Q_c - \frac{1}{2}Q_d$

[Ibanez, Marchesano, Rabadan; Cremades, Ibanez, Marchesano;'01]



Explicit realization of this structure e.g. in toroidal models

N_{lpha}	(n^1_lpha,m^1_lpha)	(n_{lpha}^2,m_{lpha}^2)	(n_{lpha}^3,m_{lpha}^3)
$N_a = 3$	(1,0)	(1,3)	(1,-3)
$N_b = 1$	(0,1)	(1,0)	(0,-1)
$N_c = 1$	(0,1)	(0,-I)	(1,0)
$N_d = 1$	(1,0)	(1,3)	(1,-3)

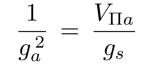
Some phenomenological properties

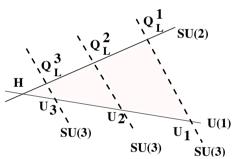
[Aldazabal, Franco, Ibanez, Rabadan, AU; ...]

- Gauge couplings
 - A priori no natural unification at string scale
 - Each coupling depends on wrapped volume
 - Can try GUT model building, see later
- 🖗 Yukawa couplings

Mediated by open string worldsheet

instantons $Y_{jk} \simeq e^{-A_{Hjk} + i\phi_{jk}}$

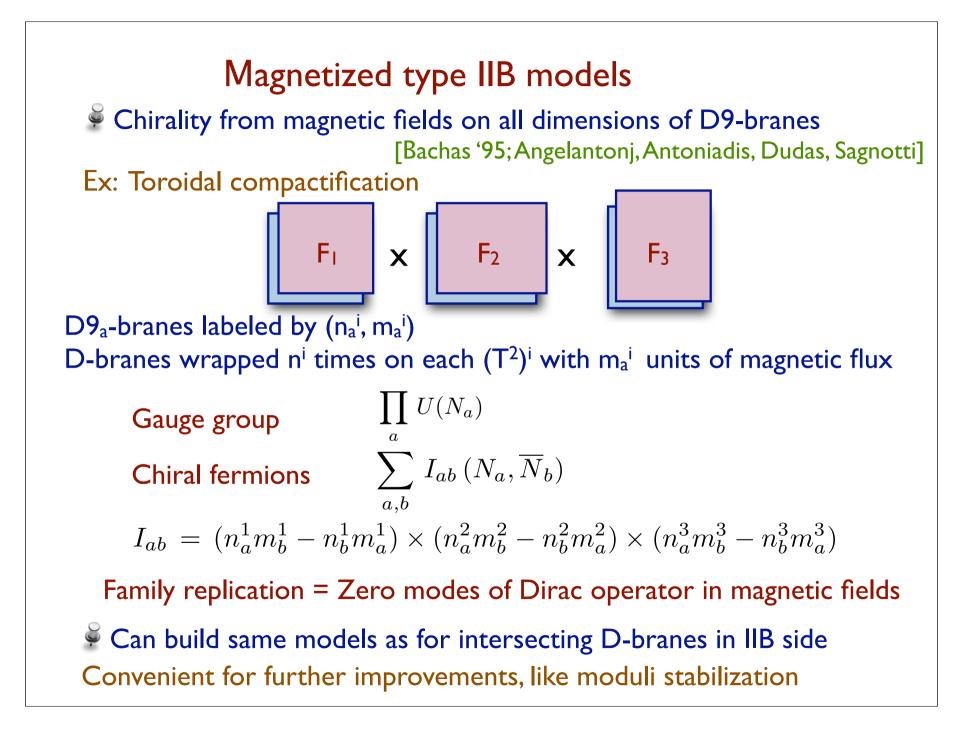




Exponential dependence potentially explains fermion mass hierarchy String scale

- Susy models, can 'choose' string scale (until susy breaking is specified)
- Non-susy models, can alleviate hierarchy by large volume [ADD'98]
- 🟺 Proton decay

In SM models, perturbatively forbidden by $U(1)_a$ baryon number Violated by instantons, just like in SM



Some phenomenological properties

[Aldazabal, Franco, Ibanez, Rabadan, AU; ...]

Gauge couplings Unified higher-dimensional gauge group, broken by magnetic field Unification modulo threshold corrections

Can try GUT model building, see later

Yukawa couplings Arise from purely field theory mechanism Overlap of zero mode wavefunction overlap

$$Y_{jk} \simeq e^{-A_{Hjk} + i\phi_{jk}}$$

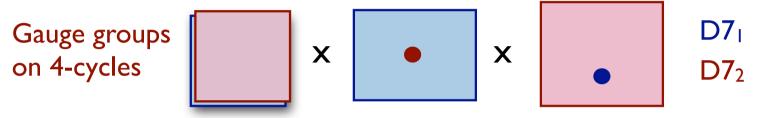
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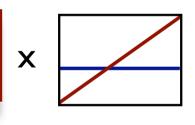
D7-brane models in type IIB

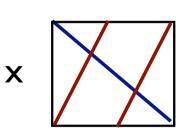
- Solution Class includes configurations with intersecting D7s with magnetic flux
 - Essentially, meaning of entries (n,m)=(0,1)



- Each D7 wraps a 4-cycle, two D7's intersect over 2-cycles

Chiral matter on 2-cycles





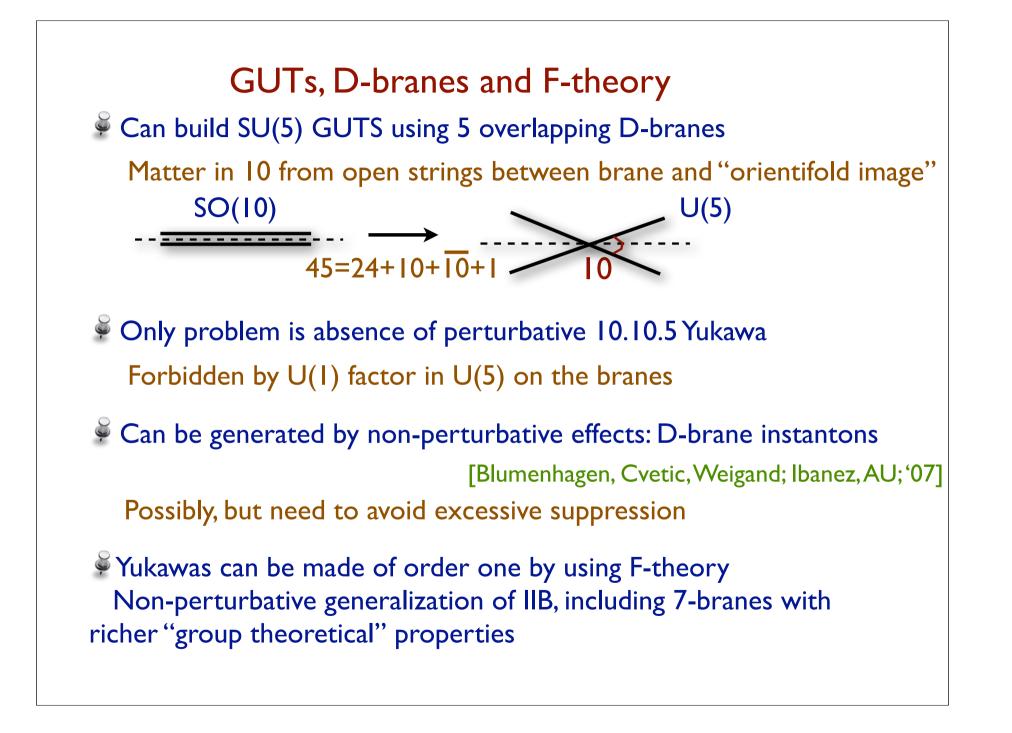
D7₂

 $D7_1$

D7₃

- Yukawa couplings localize on points

Only place where wavefunctions overlap



Yukawas in F-theory GUTS [Beasley, Heckman, Vafa]

D7₃

Solution Final Price In IIB D7's, I2-23-31 Yukawa from local enhanced U(N1+N2+N3) U(N1+N2+N3) → U(N1) × U(N2) × U(N3) Adj → Adj1 + Adj2 + Adj3 + (N1, N2) + (N2, N3) + (N1, N3) + cc. D71 D72 D71 D72

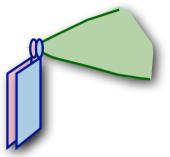
With orientifolds, 10.5.5 Yukawa from local SO(12) SO(12) \rightarrow SO(10) \times U(1) \rightarrow SU(5) \times U(1) \times U(1) 66 \rightarrow 45 + 10 + 10 + 1 \rightarrow 24 + 10 + 5 + 5 + cc + singlets

In F-theory, 7-branes with local E6 enhancement produce 10.10.5 E6 → SO(10) × U(1) → SU(5) × U(1) × U(1) 78 → 45 + 16 + 16 + 1 → 24 + 10 + 10 + 5 + cc + singlets

F-theory allows to reconcile brane model building and GUT ideas

Local F-theory models

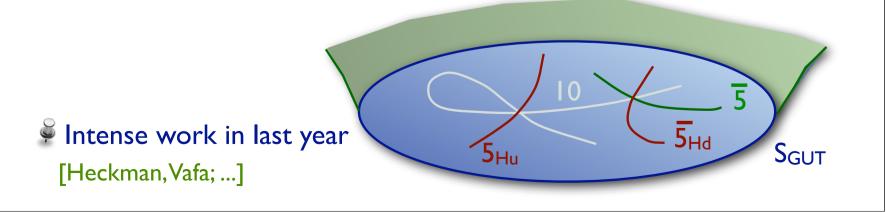
- Bottom-up approach [Aldazabal, Ibanez, Quevedo, AU; '00]
 Build local configuration of D-branes, later embed in compact space
- 🗳 Ingredients
 - SU(5) 7-brane wrapped on a "small" 4-cycle S_{GUT}
 - " del Pezzo surface": Blowup of a singularity
 - U(I) hypercharge magnetic flux breaks to SM



[Donagi,Wijnholt;

Beasley, Heckman, Vafa]

- Other 7-branes intersect S_{GUT} over 2-cycles, leading to chiral matter SM fermions from 2-cycles with no hypercharge flux: SU(5) multiplets Higgses from 2-cycles with hypercharge flux: DT splitting
- -Yukawas from points where matter 2-cycles intersect



Global GUT models

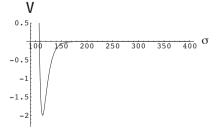
- ✓ In type IIB orientifolds [Blumenhagen, Braun, Grimm, Weigand, '08;]
 Type IIB analog of F-theory GUTS
 Build 6d CY compact space with a local del Pezzo 4-cycle
 Wrap D7-branes to reproduce local structure of GUT model
 Satisfy certain global consistency conditions (Gauss law)
 Require existence of non-perturbative instantons generating 10.10.5
 ⇒ Explicit model with 3 families
- 🗳 In F-theory

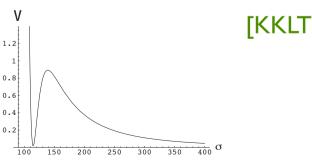
[Marsano, Saulina, Schafer-Nameki, '09]

7-branes source complex coupling T, which varies over 6d space Geometrize as 2 extra T² dimensions \Rightarrow 8d compact space 7-brane magnetic field becomes 4-form flux in this 8d space Build 8d compact CY space with local patch reproducing F-theory GUT Suceed in building the geometry, but flux is still difficult to handle \Rightarrow On the way to explicit model with 3 families

Moduli stabilization and fluxes

- Simplest compatifications ⇒ moduli, massless scalars w/ no potential
 e.g. parametrize sizes of 2- and 3-cycles in CY (Kahler and complex)
- Moduli stabilization from fluxes in the compactification Field strength fluxes on cycles, geometric fluxes, ... Nontrivial dependence on moduli => scalar potential [Dasgupta, Rajesh, Sethi; Giddings, Kachru, Polchinski; ...]
- Explicit models of (MS)SM with (partial) flux moduli stabilization
 [Blumenhagen, Lust, Taylor; Cascales, AU; Camara, Font, Ibanez; Villadoro, Zwirner;
 '03-'06]
 - Moduli stabilization & susy breaking in Minkoswki (full stab. in AdS) Going to dS is open question, with interesting but not yet explicit proposals





Fluxes, susy breaking and soft terms

- Appealing (not unique) scenario: Susy MSSM D-branes and non-susy flux
- Soft terms arise from effect of non-susy flux on susy D-branes

Explicitly computable using D-brane world-volume action in general supergravity background, or using 4d effective theory approach [Grana; Camara, Ibanez, AU; Lust, Mayr, Reffert, Stieberger; '03-'04]

- $\frac{1}{2}$ Flux components work as vevs for auxiliary fields of chiral multiplets of (complex structure) moduli
 - \Rightarrow Realization of gravity-mediated susy breaking
 - Flavour problem: Decoupling of flavor physics and soft terms Geometrization squark masses determined by intersection angles
 - μ -problem: susy components of flux induce it on the branes
 - Very explicit discussion of susy spectrum etc is possible in specific models

e.g. in 'large volume compactifications' [Quevedo et al '06-'07] F-theory [Aparicio, G.Cerdeño, Ibáñez, '08]

Conclusions

- $\frac{1}{2}$ It is remarkable that there are string models so close to (MS)SM
- 🗳 Unique theory, yet many vacua

Many choices of data of compactification to 4d

Fules are very well defined

Many examples of things which are just not possible

- -Very large representations \Rightarrow e.g. no 126 of SO(10)
- Bounds on number of branes ⇒ rank of gauge groups
 Many examples of things not possible in a particular class
 No SO(10) GUT on weakly coupled D-brane models

Role of geometric intuition: Geometrization of SM structures

Role of modular structures:

Relatively independent building blocks for SM, susy breaking, inflation, moduli stabilization, ...

What is it good for?

Many realistic vacua: No unique testable prediction
 Each particular consistent realistic model is probably wrong
 But some general lessons may be right and key to the UV of SM

- New scenarios (in UV complete theory): Extra dimensions, brane world, warping, ...

Plausible patterns within each
 e.g. Low energy susy and susy breaking soft terms

- Smoking guns for some scenarios (±contrived) e.g. string resonances in TeV scale models

- Impact of LHC results?

Expect continuous progress in understanding the UV of SM