

The cosmic ray antiproton signal at TeV energies

Pierre Salati – Université de Savoie & **LAPTH**

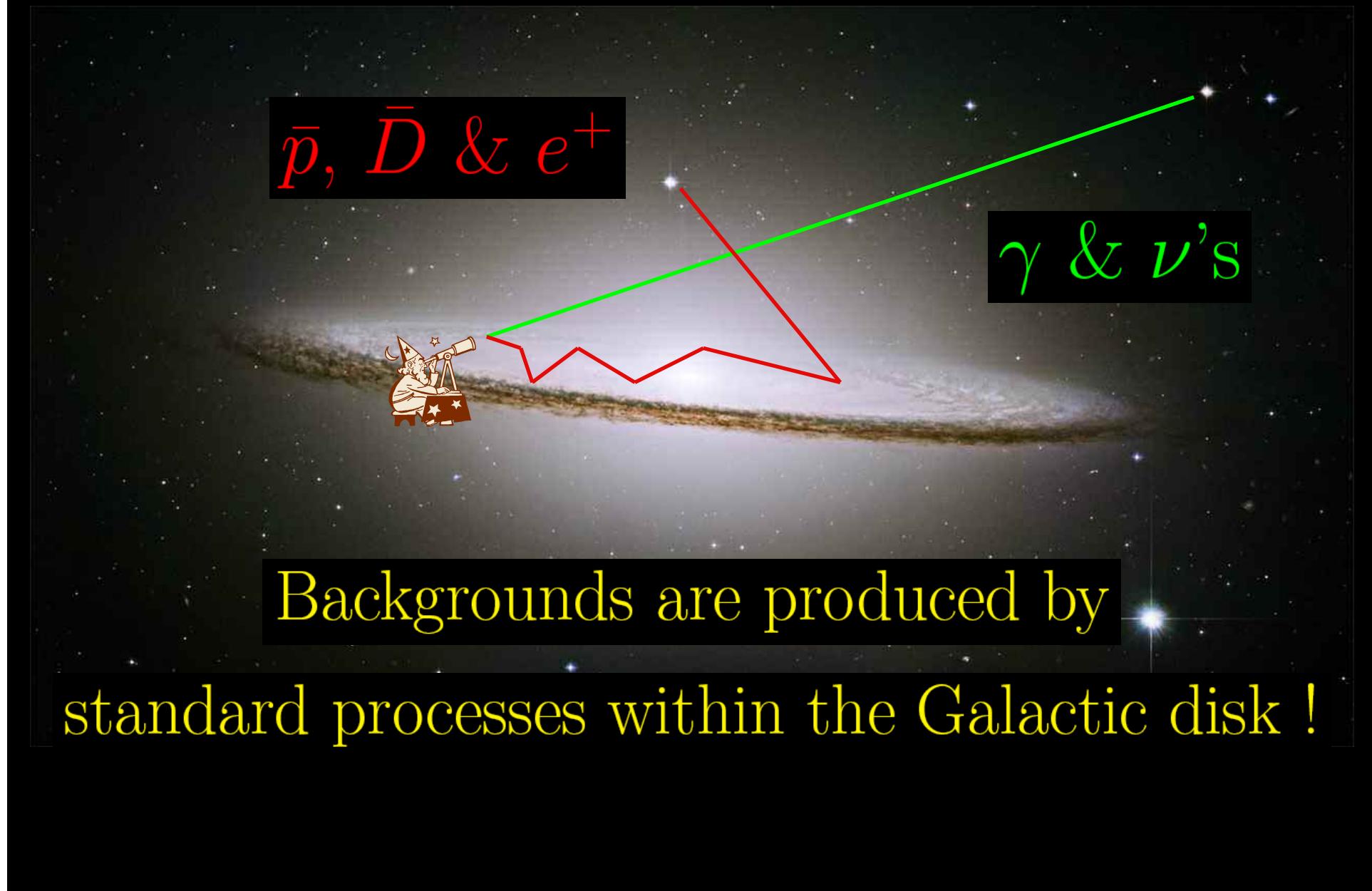
- 1) High energy antiproton : a new window
- 2) The antiproton signal of TeV WIMPs
- 3) DM mini-spikes around IMBHs

T. Bringmann & P.S., PRD **75** (2007) 083006

P. Brun, G. Bertone, J. Lavalle, P.S. & R. Taillet, arXiv:0704.2543

TeV Particle Astrophysics 2007 – Venezia – August 28, 2007

1) High energy antiproton : a new window



Cosmic-rays diffuse in space and energy

- Diffusion and convection in space

$$\vec{J} = -K \vec{\nabla} \Psi + \Psi \vec{V}_C$$

- Second order Fermi mechanism

$$J_E = b^{\text{loss}}(E) \Psi - K_{EE}(E) \partial_E \Psi$$

- Steady state holds with $\partial_t \Psi = 0$ Thickness L

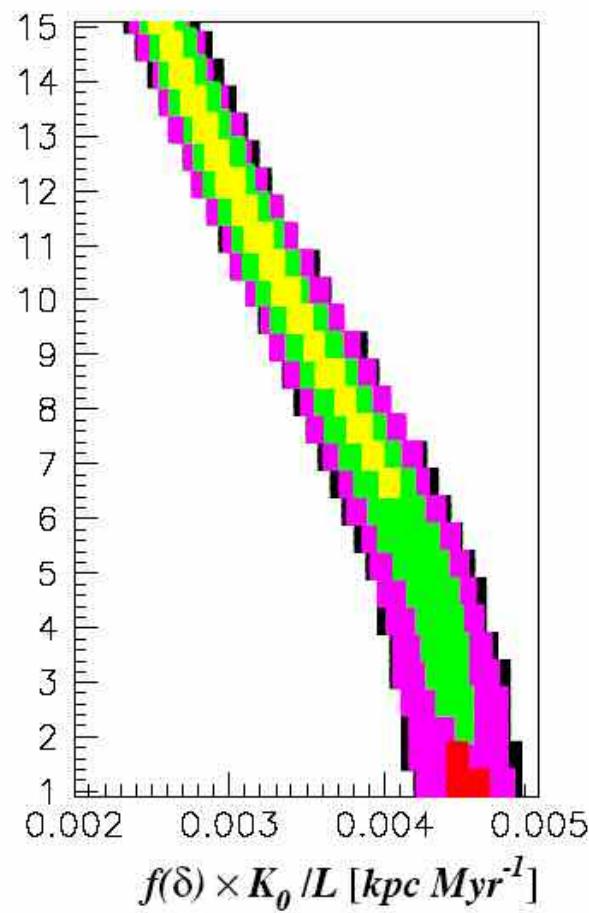
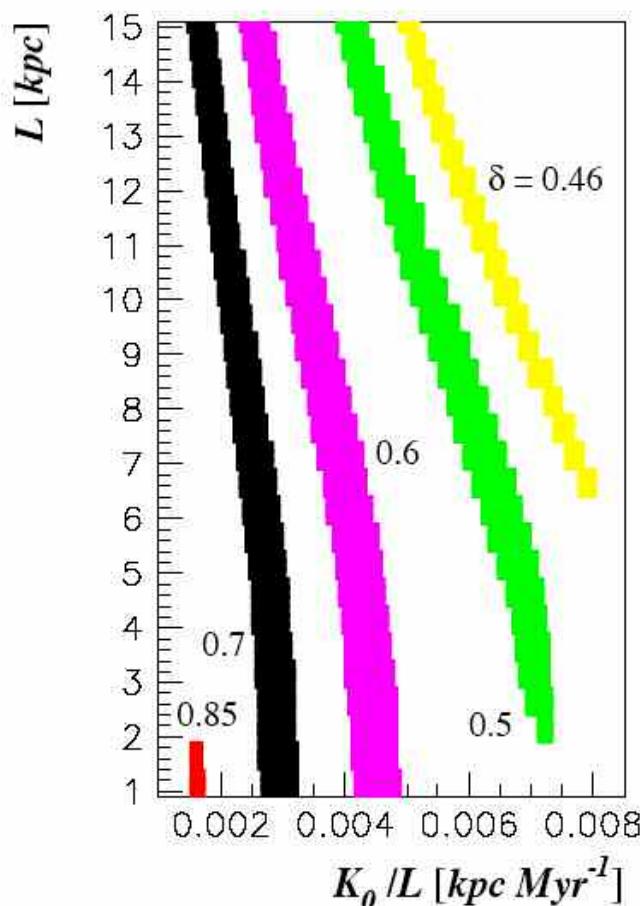
$$K(E) = K_0 \beta \times \mathcal{R}^\delta \quad \Downarrow \quad K_{EE} = \frac{2}{9} V_a^2 \frac{E^2 \beta^4}{K(E)}$$

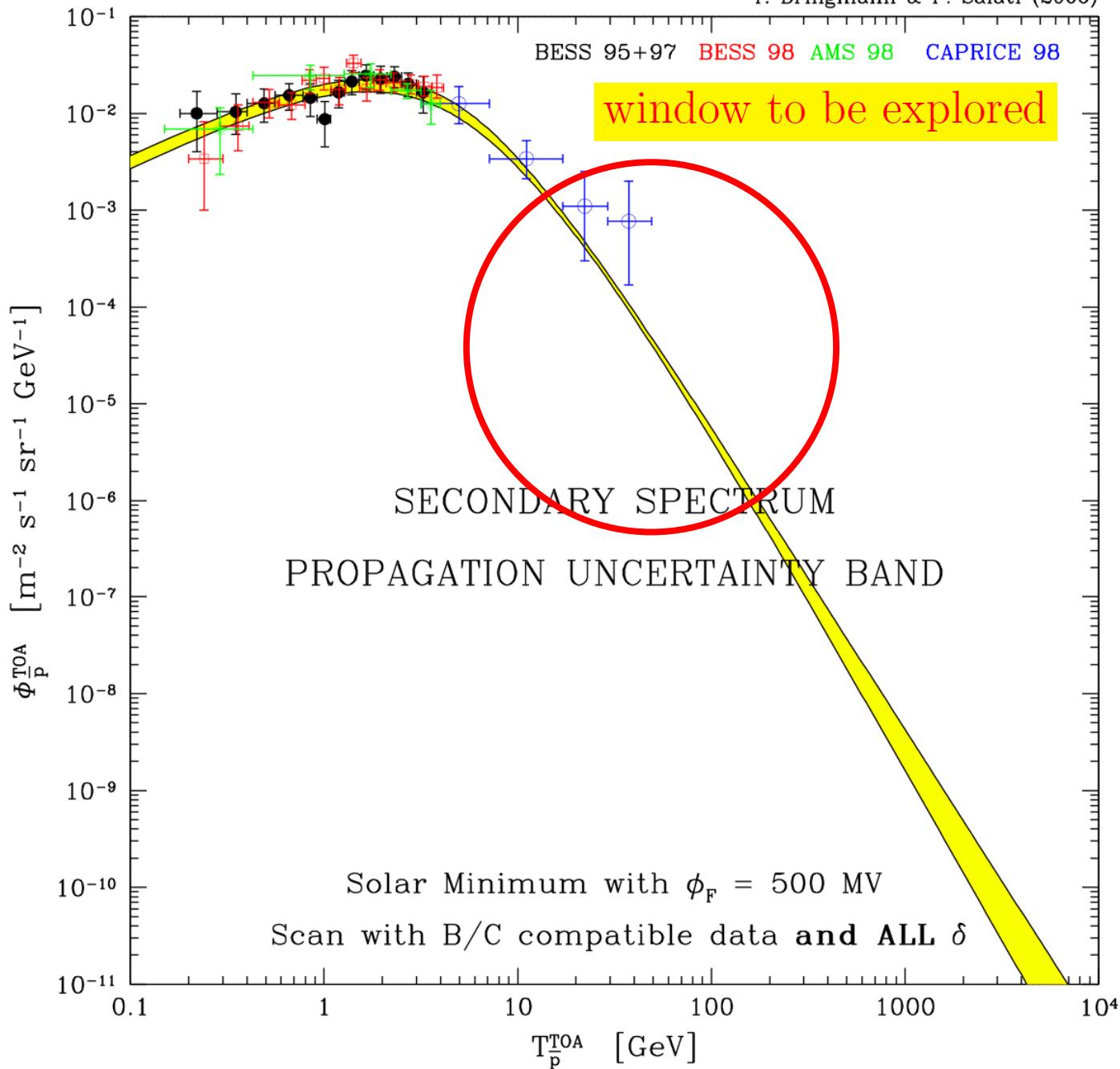
$$V_C \partial_z \Psi - K \Delta \Psi + \partial_E \{ b^{\text{loss}}(E) \Psi - K_{EE}(E) \partial_E \Psi \} = Q$$

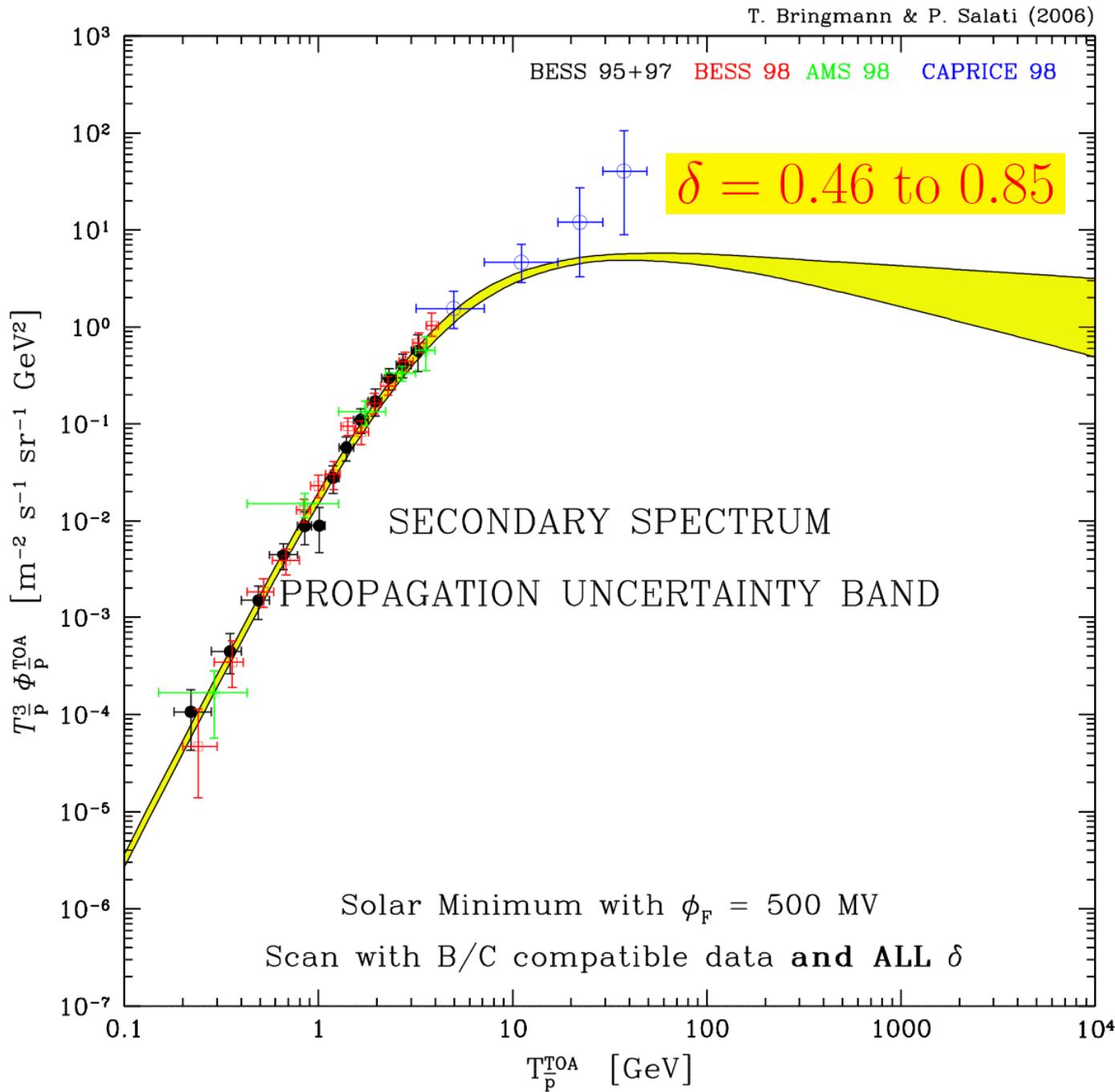
D. Maurin et al.

$$K(E) = K_0 \beta \times \mathcal{R}^\delta$$

Iso- χ^2 contours for B/C ($\chi^2 < 40$)







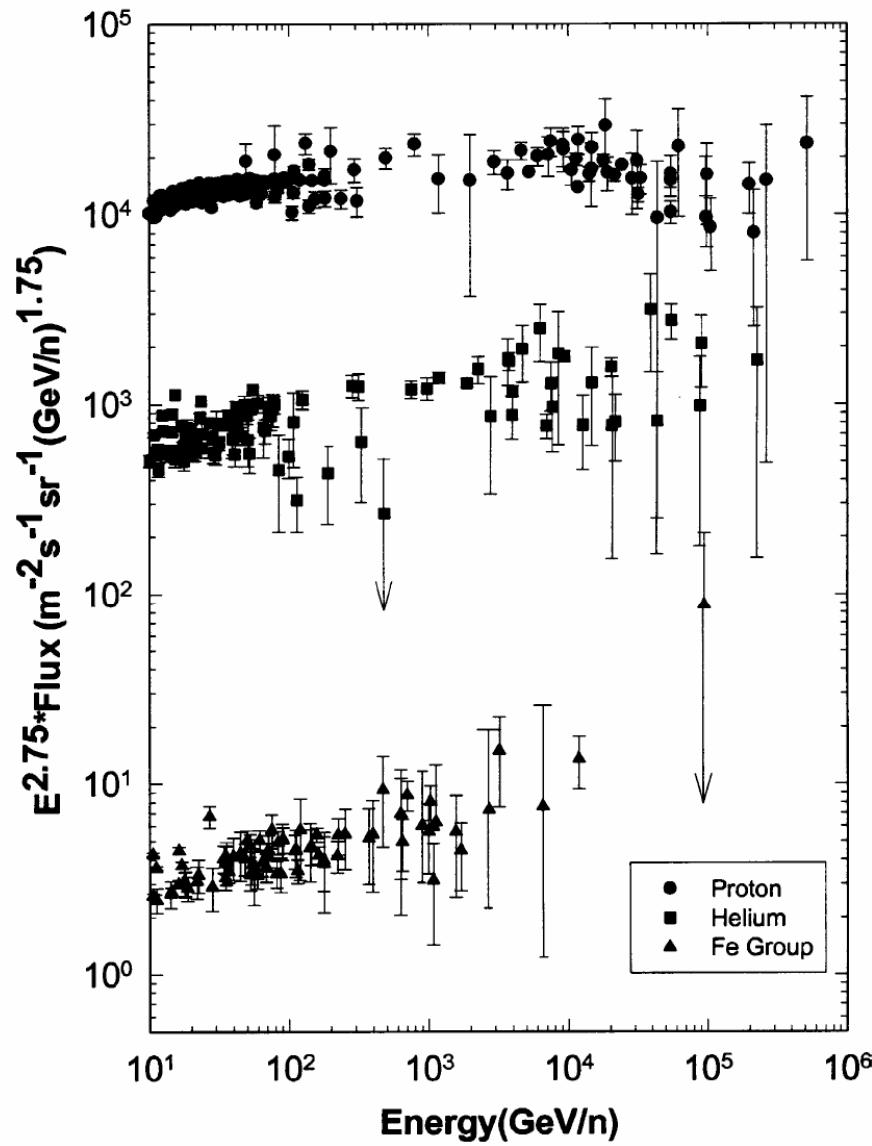
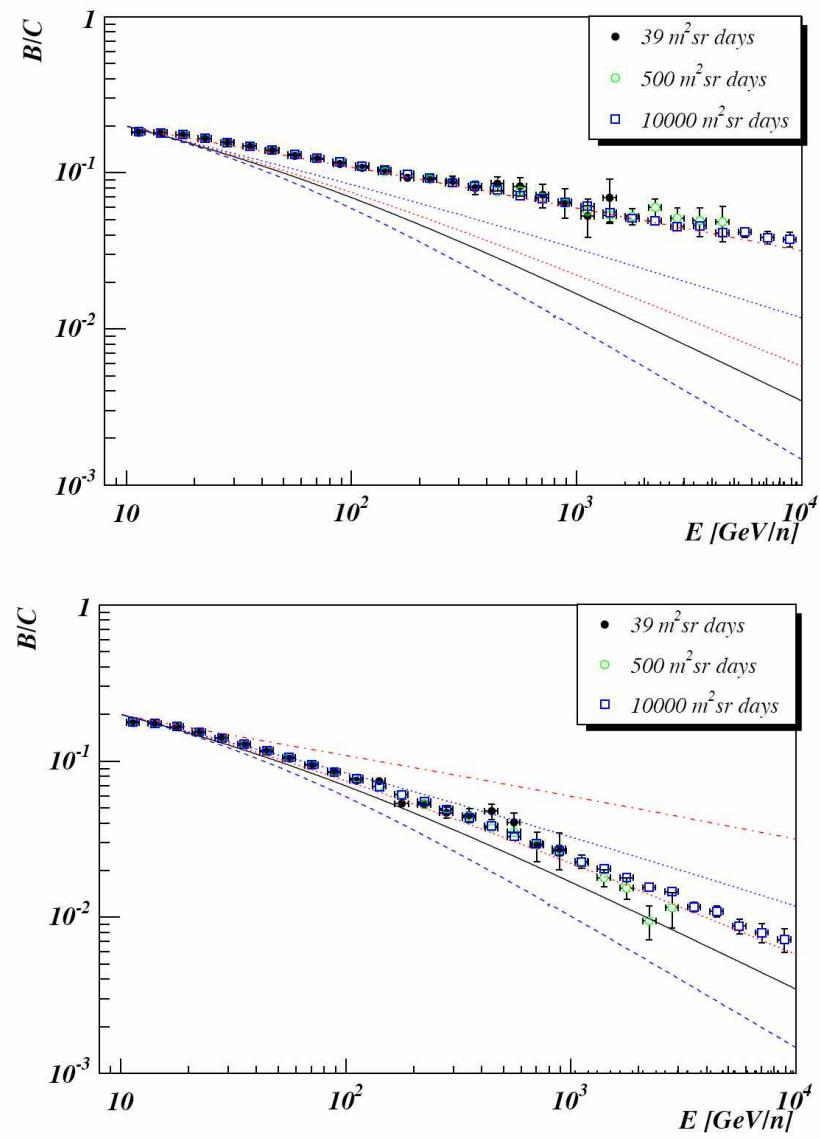


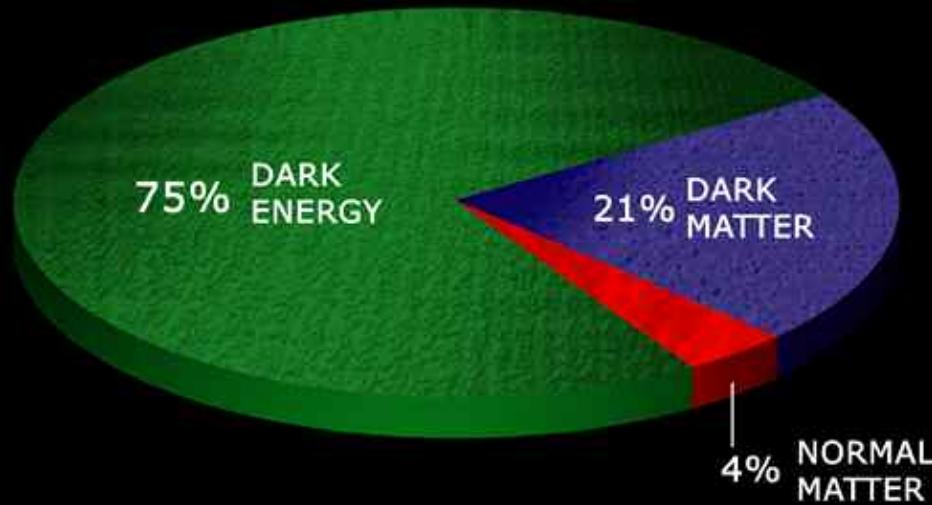
Fig. 1. Compiled data from direct measurements.

CREAM



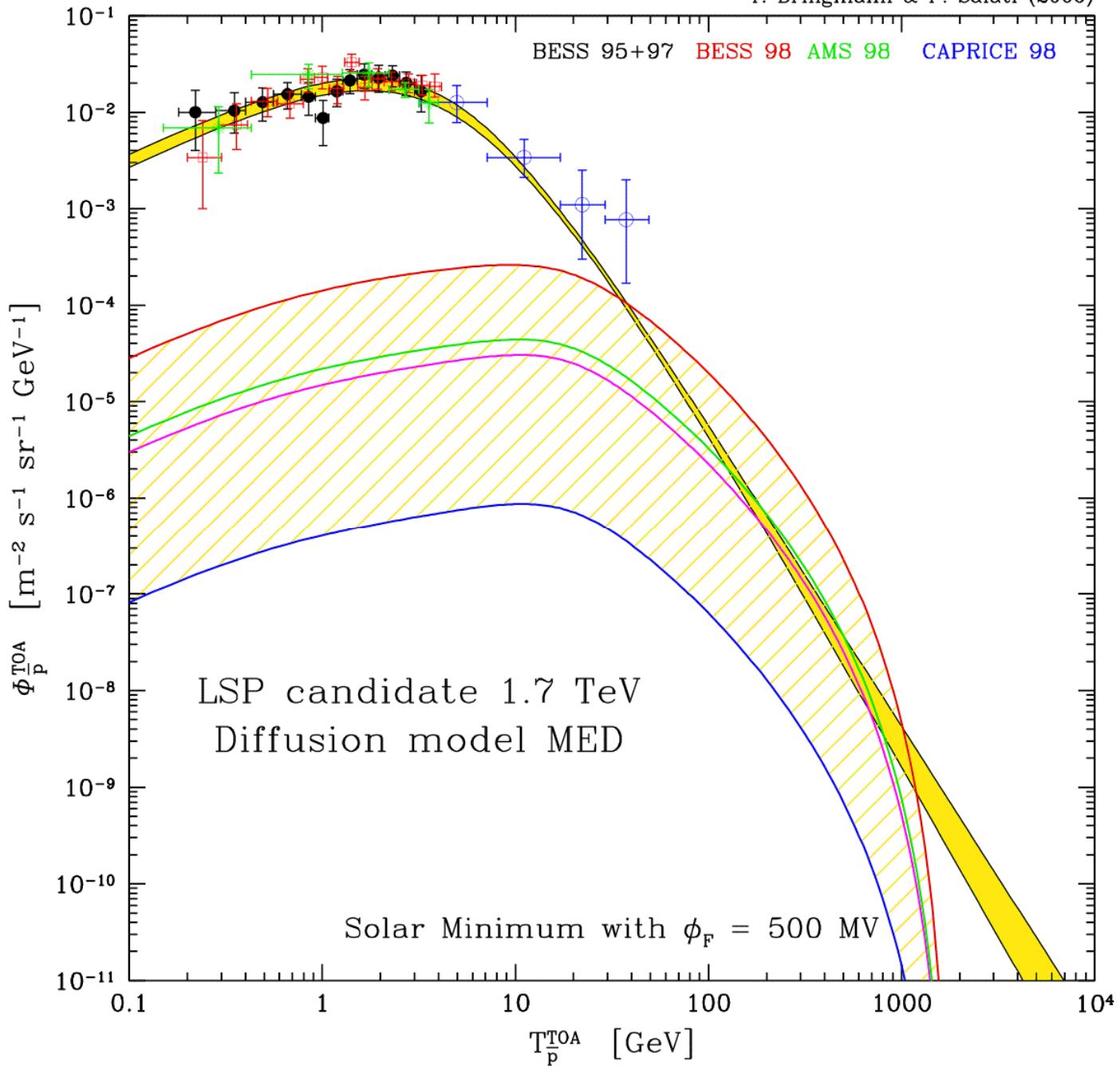
A. Castellina & F. Donato astro-ph/0504149

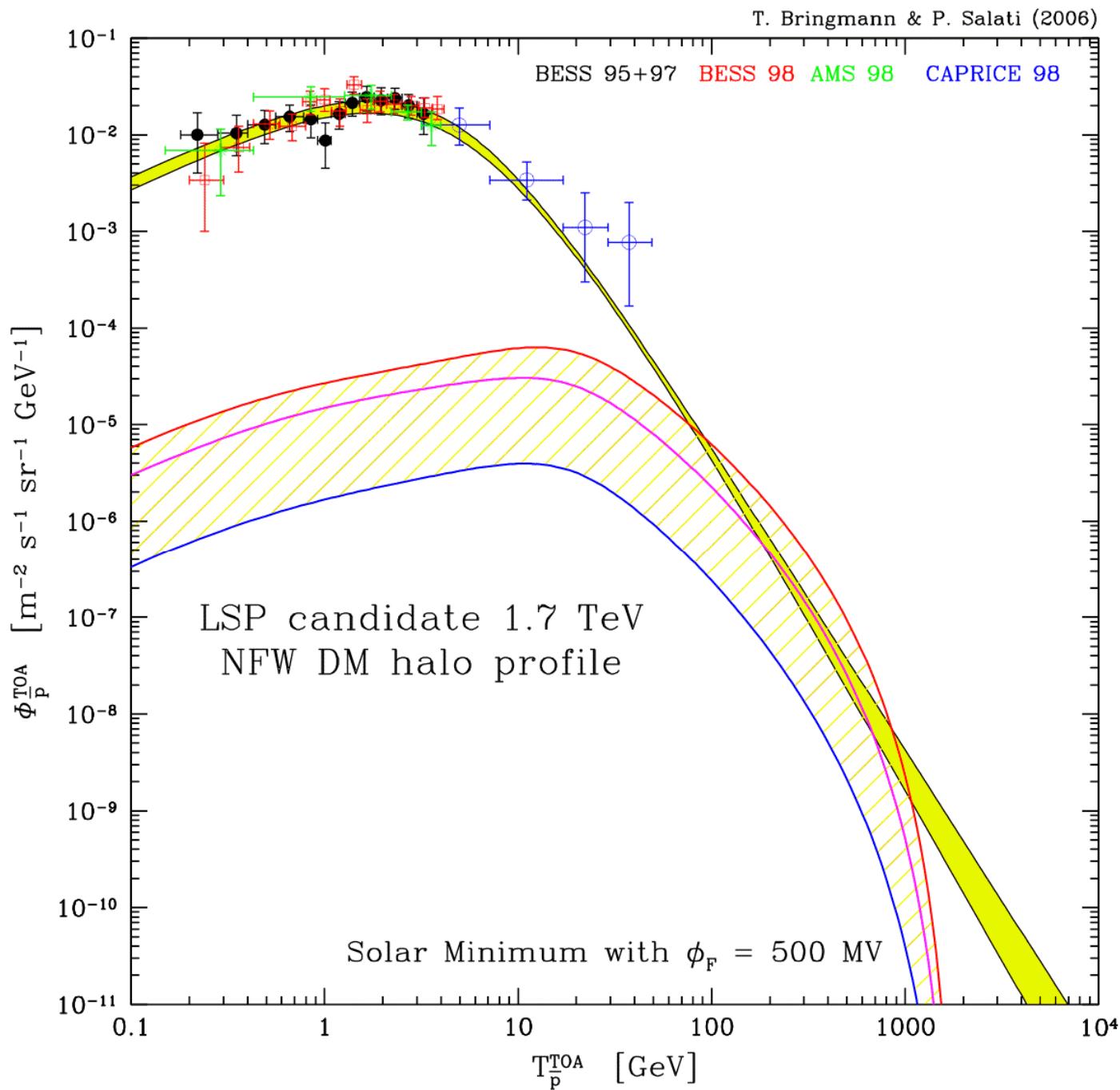
2) The antiproton signal of TeV WIMPs

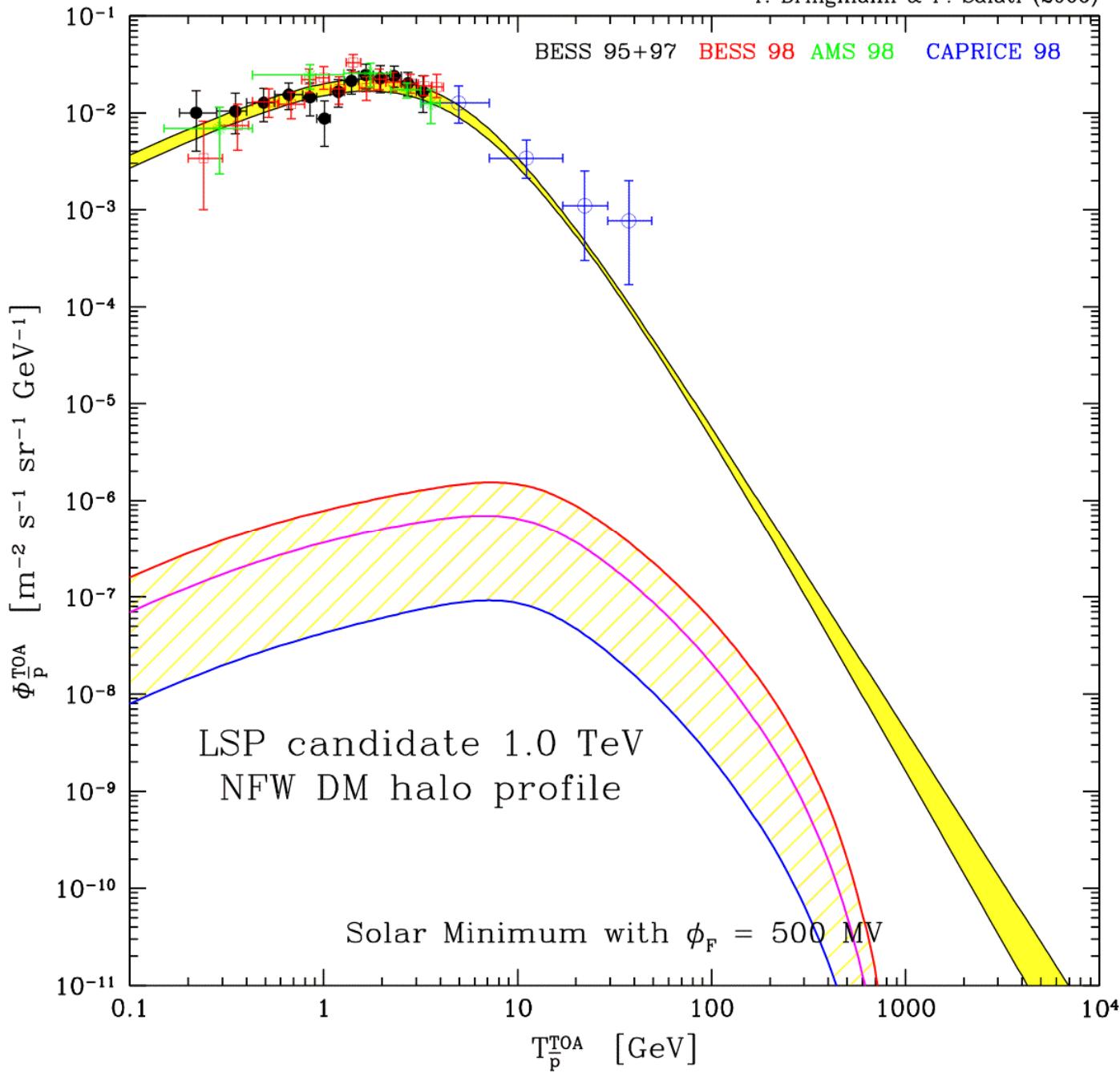


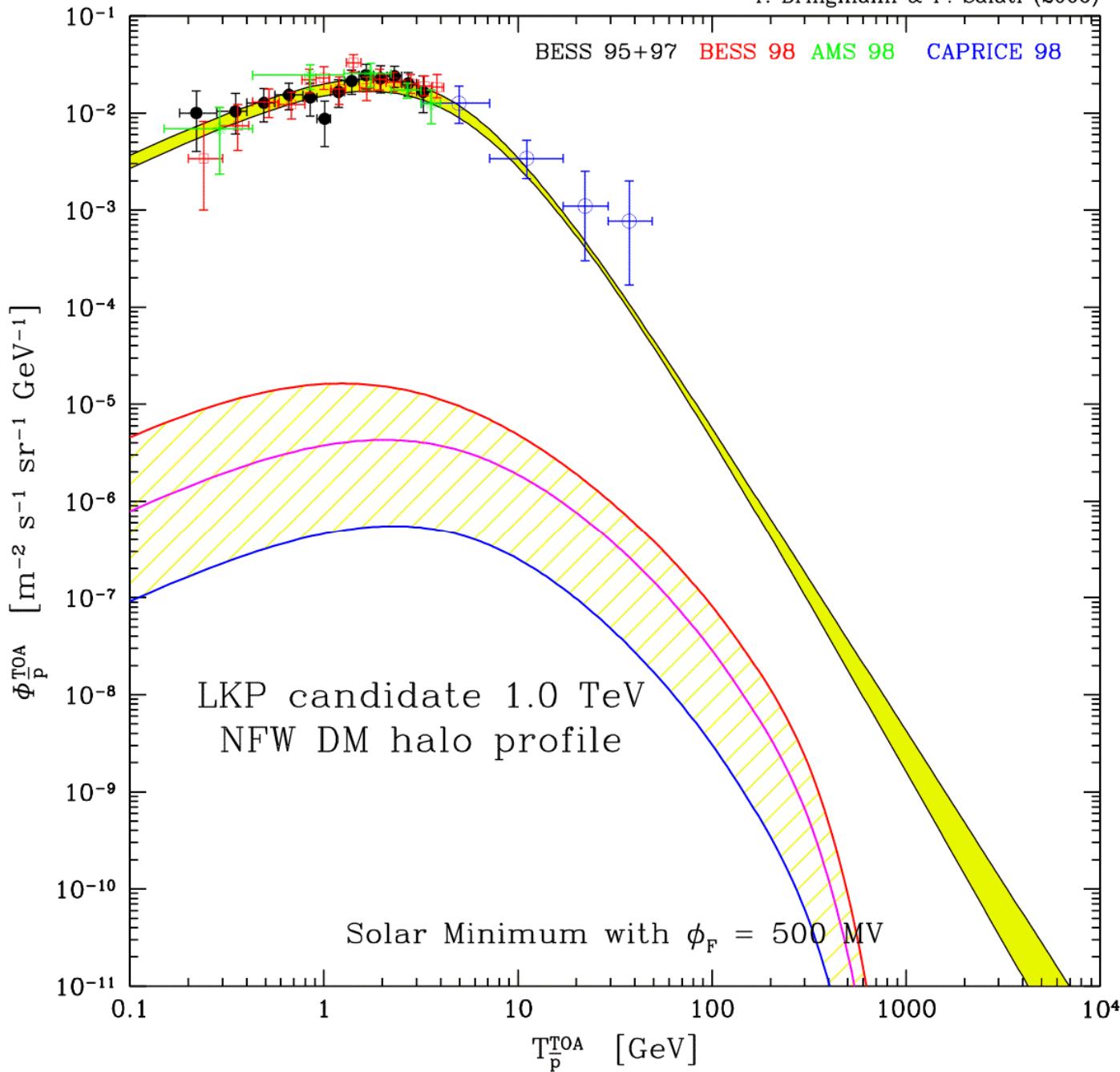
$$\chi + \chi \rightarrow f\bar{f}, W^+W^-, \dots \rightarrow \bar{p}, \bar{D}, e^+, \gamma \text{ & } \nu's$$

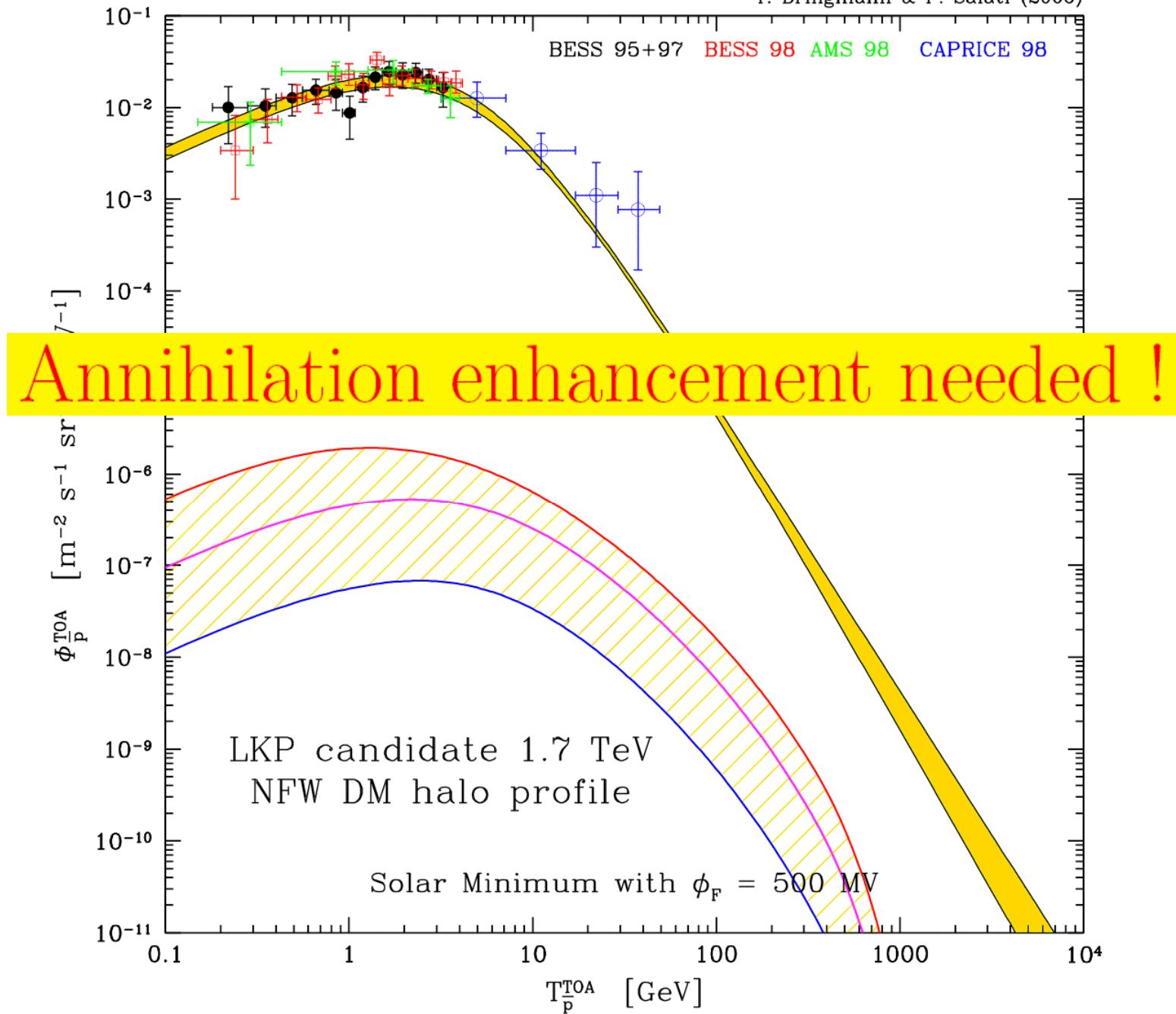
Weakly Interacting Massive particles – WIMPs – could be the major component of the haloes of galaxies. Their mutual annihilations would produce extra high-energy cosmic rays : the DM messengers

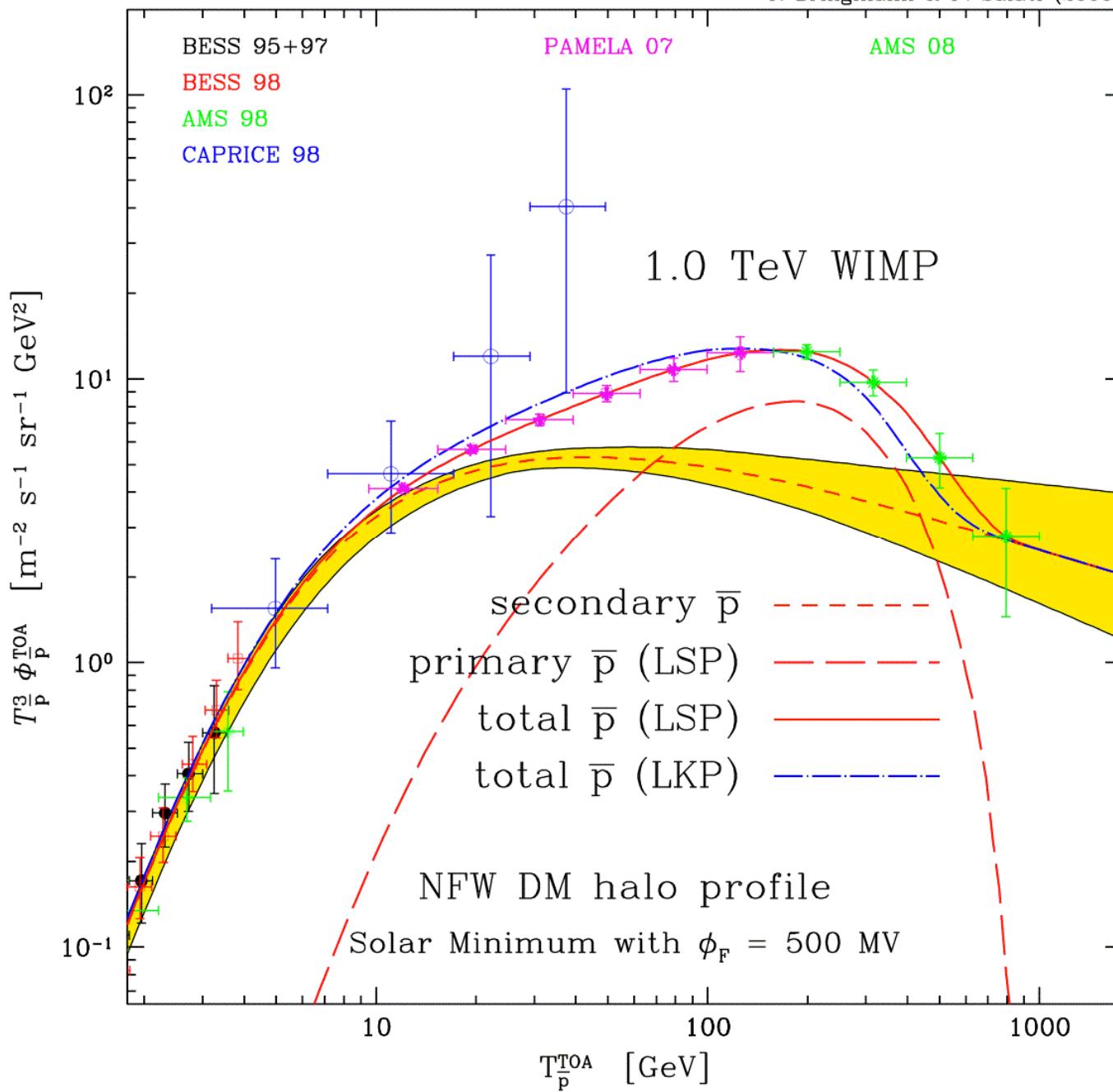


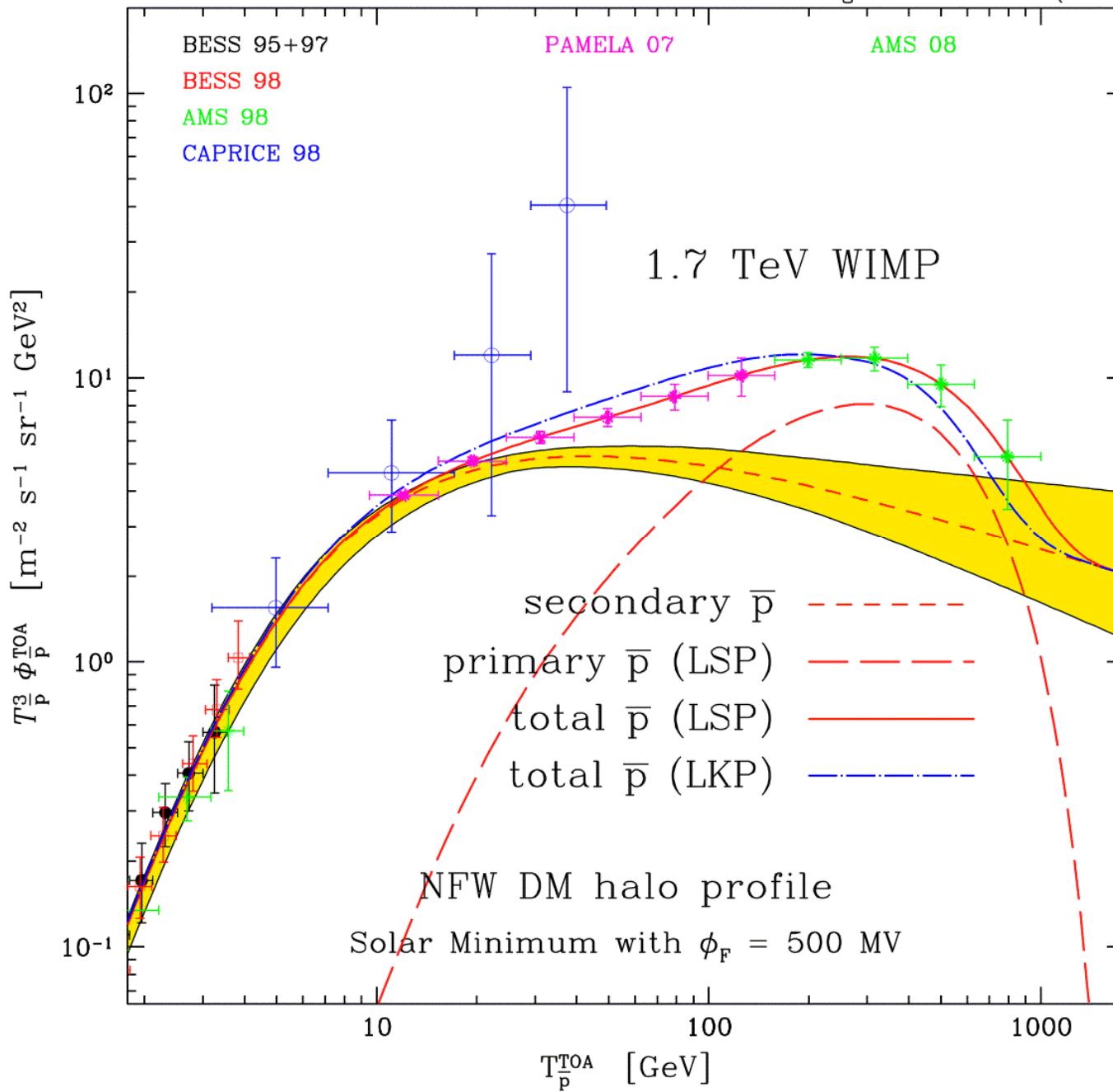












3) DM mini-spikes around intermediate mass black holes

G. Bertone, A.R. Zentner & J. Silk, PRD **72** (2005) 103517

When the first DM halos form, gas cools and collapses as pressure supported disks

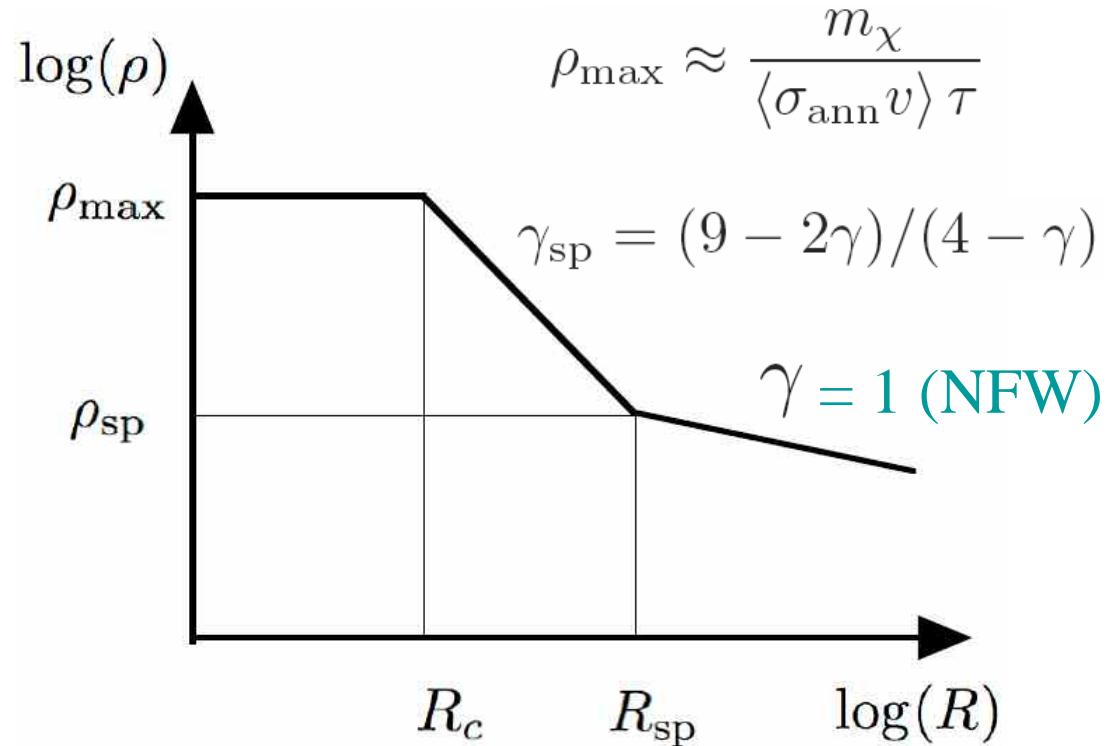
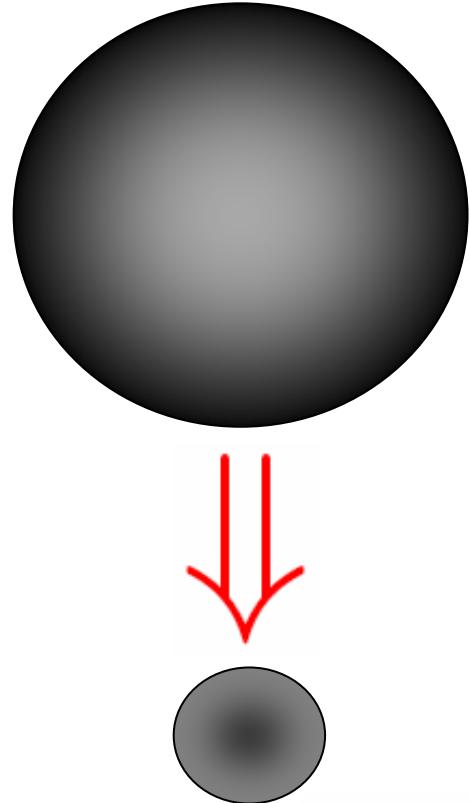


A baryonic mass of $\sim 10^5 M_\odot$ loses its angular momentum

It is transferred at the center to form an Intermediate Mass Black Hole

During the process, DM is adiabatically compressed onto this central object

Adiabatic DM compression around the IMBH



Large annihilation volume 1.84 pc^3

$$\xi = \frac{5 \pi \rho_{\text{sp}}^2 (\rho_\odot)^{9/7} R_c^{9/7}}{\langle \sigma_{\text{ann}} v \rangle^{2/7} m_\chi^{-9/7}} \equiv \frac{\rho_{\max}}{\rho_{\text{sp}}}^{3/7}$$

Compensation between ρ_{\max} & $\langle \sigma_{\text{ann}} v \rangle$ [eV]

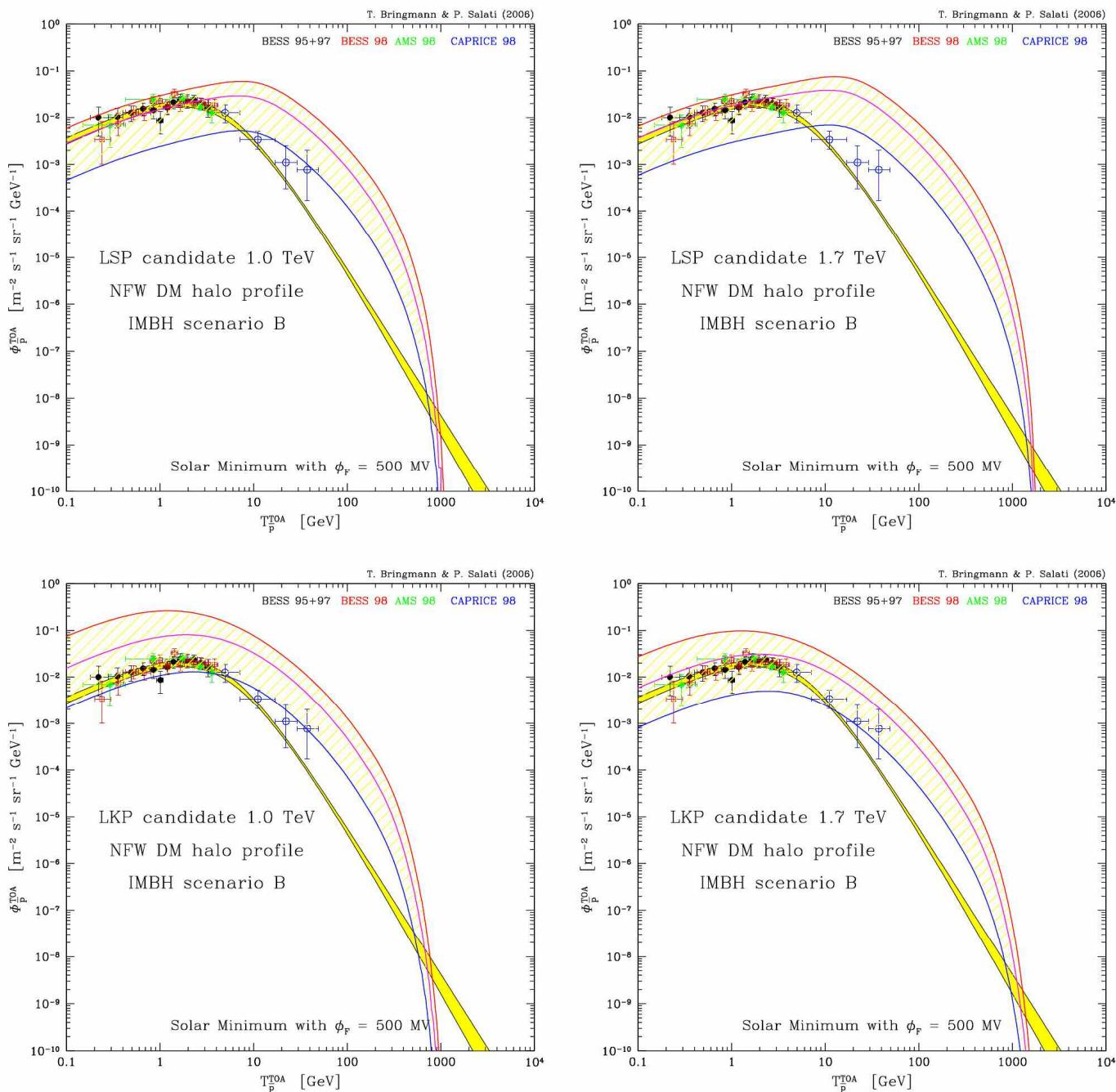


FIG. 6: Same as Fig. 5, now with an NFW profile and a typical population of IMBHs in the galactic halo.

MC distributions of the annihilation volume ξ and position \mathbf{x}

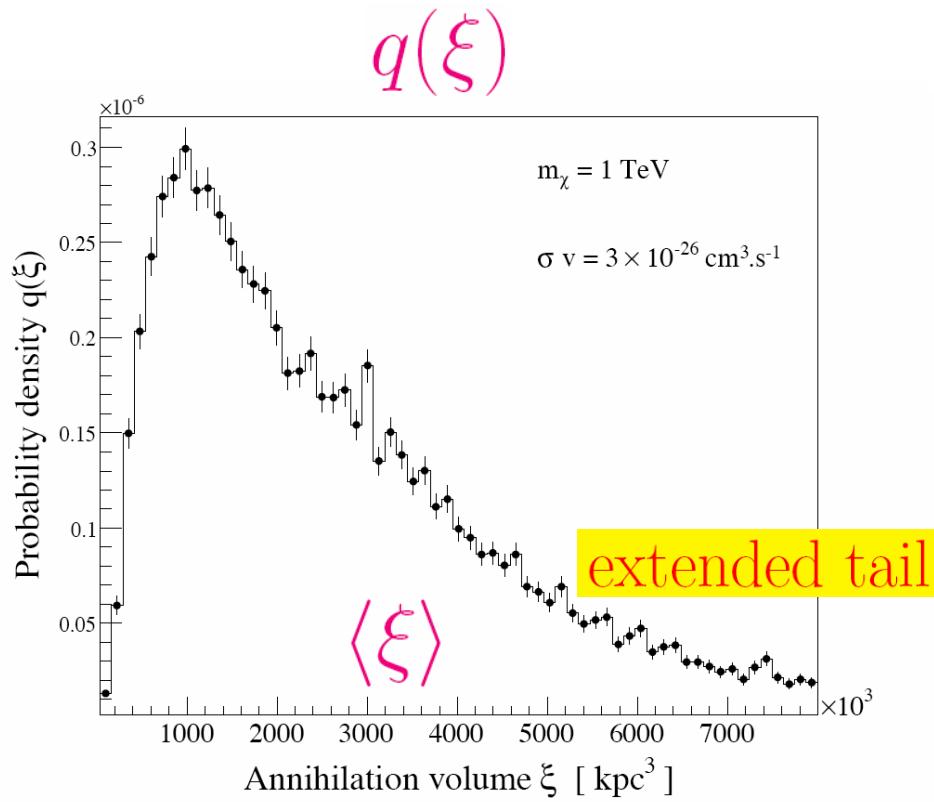


FIG. 3: The probability law $q(\xi)$ for the annihilation volume has been derived from the Monte-Carlo results of Ref. [40]. We found no correlation with the mini-spike position.

$$\langle \xi^n \rangle = \int_0^{+\infty} \xi^n q(\xi) d\xi \quad \& \quad \langle G^n \rangle = \int_{\text{DZ}} \{G(\mathbf{x})\}^n p(\mathbf{x}) d^3\mathbf{x}$$

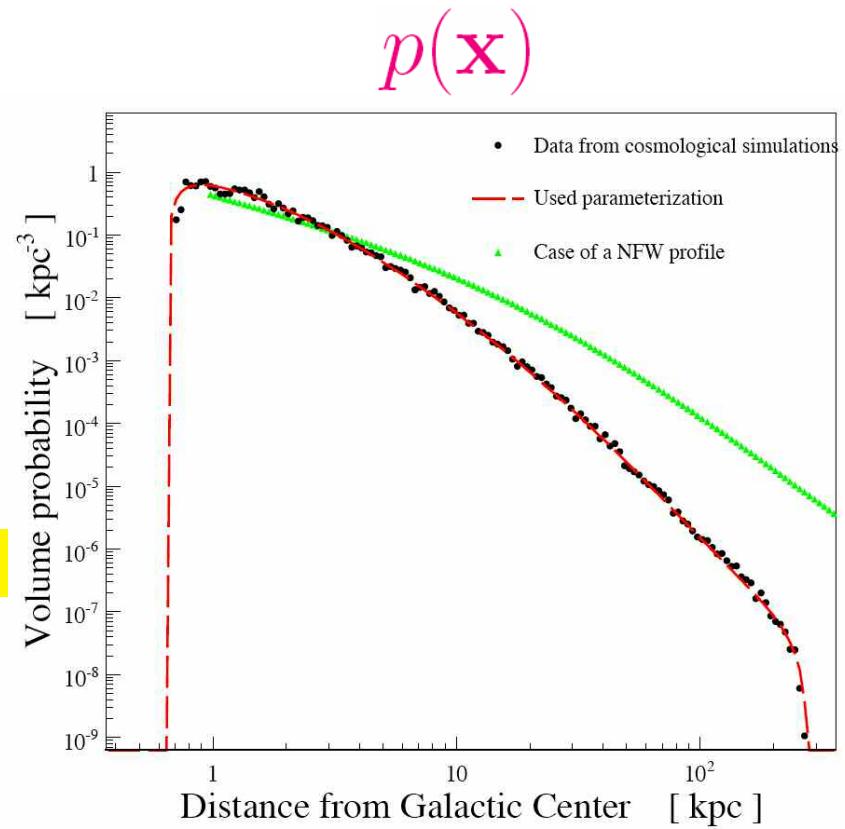
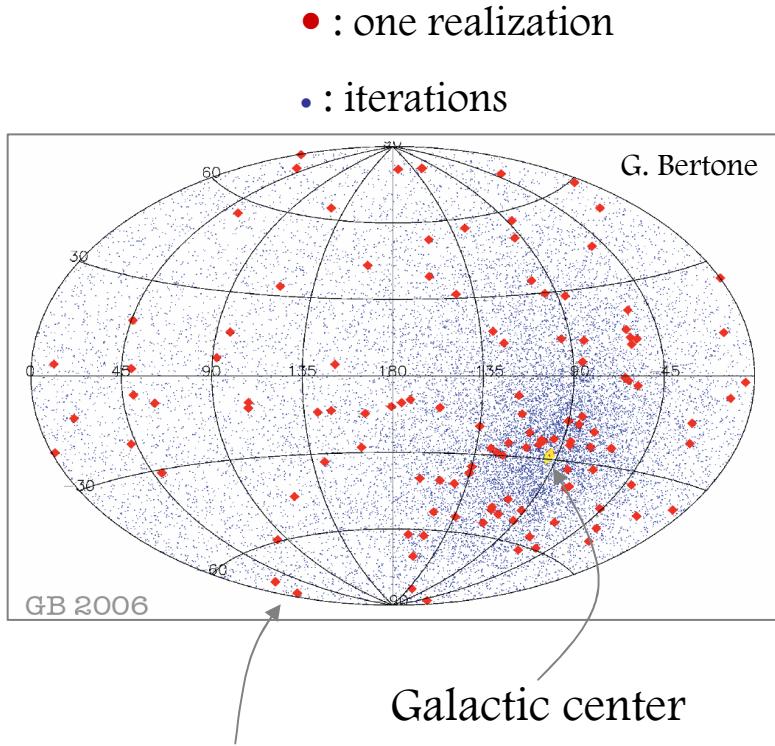


FIG. 4: Radial distribution of the mini-spikes, as extracted from the numerical results of Ref. [40].



$$B_{\text{eff}} = 1 + \frac{\langle \phi_r \rangle}{\phi_s} = 1 + N_{\text{BH}} \frac{\langle \xi \rangle \langle G \rangle}{\mathcal{I}}$$

$$\frac{\sigma_r^2}{\langle \phi_r \rangle^2} = \frac{1}{\langle N_{\text{BH}} \rangle} \left\{ \frac{\langle \xi^2 \rangle}{\langle \xi \rangle^2} \frac{\langle G^2 \rangle}{\langle G \rangle^2} - 1 \right\} + \frac{\sigma_N^2}{\langle N_{\text{BH}} \rangle^2}$$

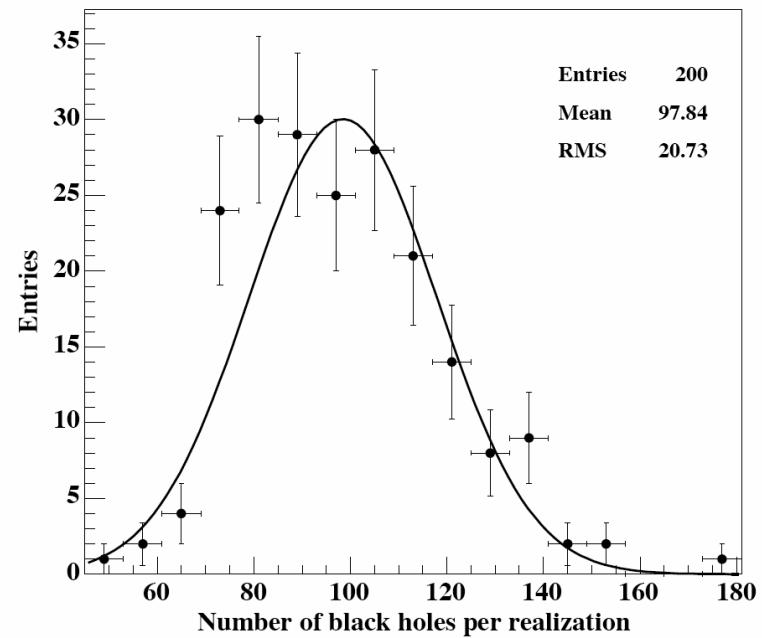


FIG. 1: Distribution of the Monte-Carlo realizations of the galactic mini-spike population – extracted from Ref. [40] – as a function of the number N_{BH} of objects within a galactocentric radius of 100 kpc.

$$\frac{\sigma_B}{B_{\text{eff}}} = \frac{\sigma_r/\phi_s}{1 + \langle \phi_r \rangle/\phi_s} \simeq \frac{\sigma_r}{\langle \phi_r \rangle}$$

Boost distributions

Positrons

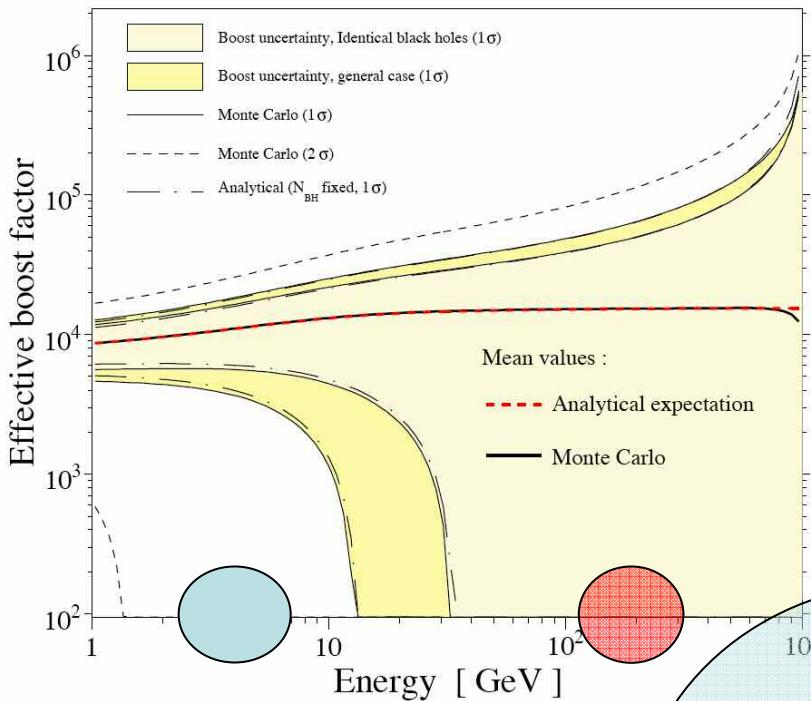


FIG. 5: Results from the Monte-Carlo simulations of the IMBHs population inside the Milky Way are compared to the analytical computations of the effective boost factor and its dispersion, for $m_\chi = 1 \text{ TeV}$.

1 TeV positron line

Antiprotons

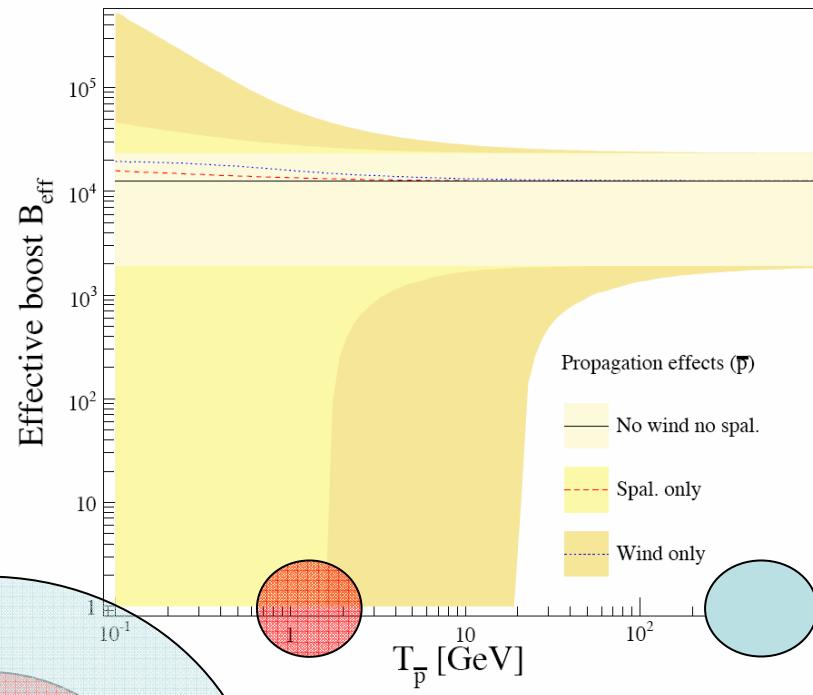
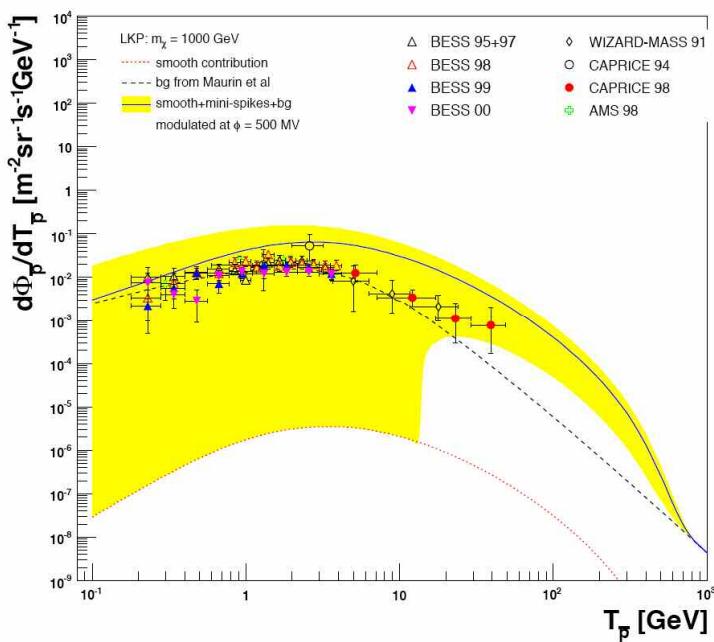
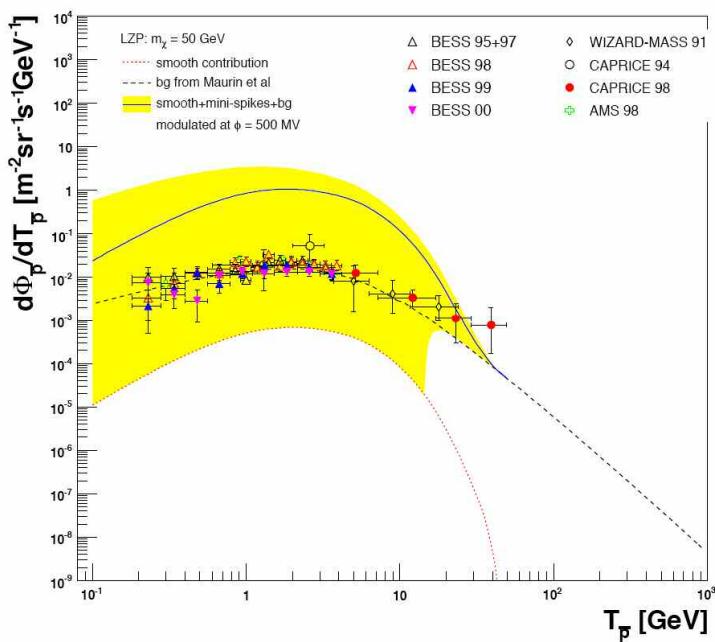
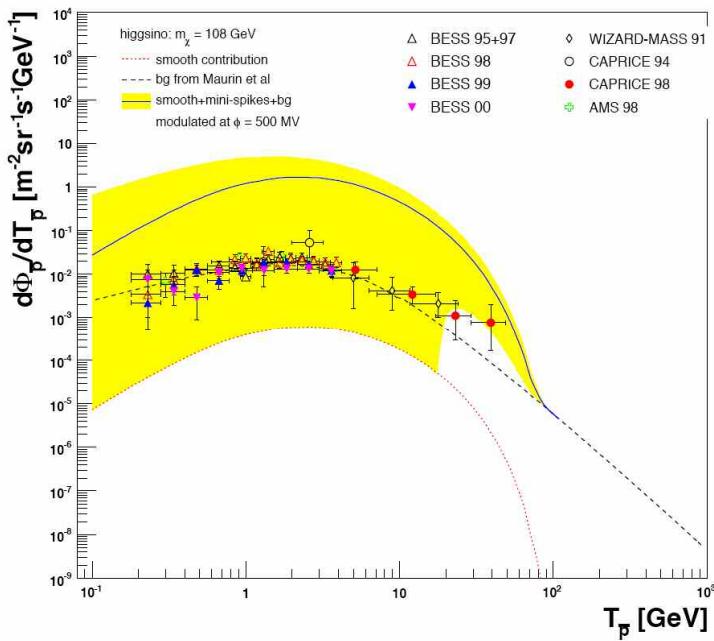
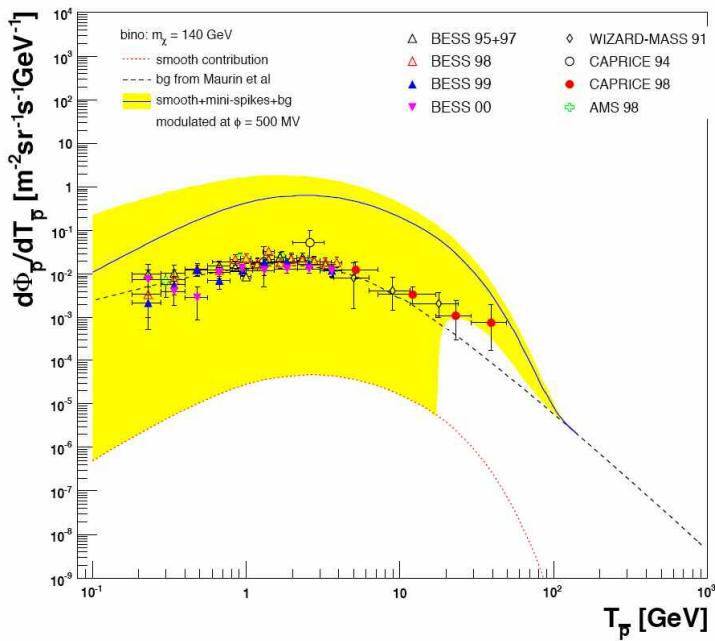


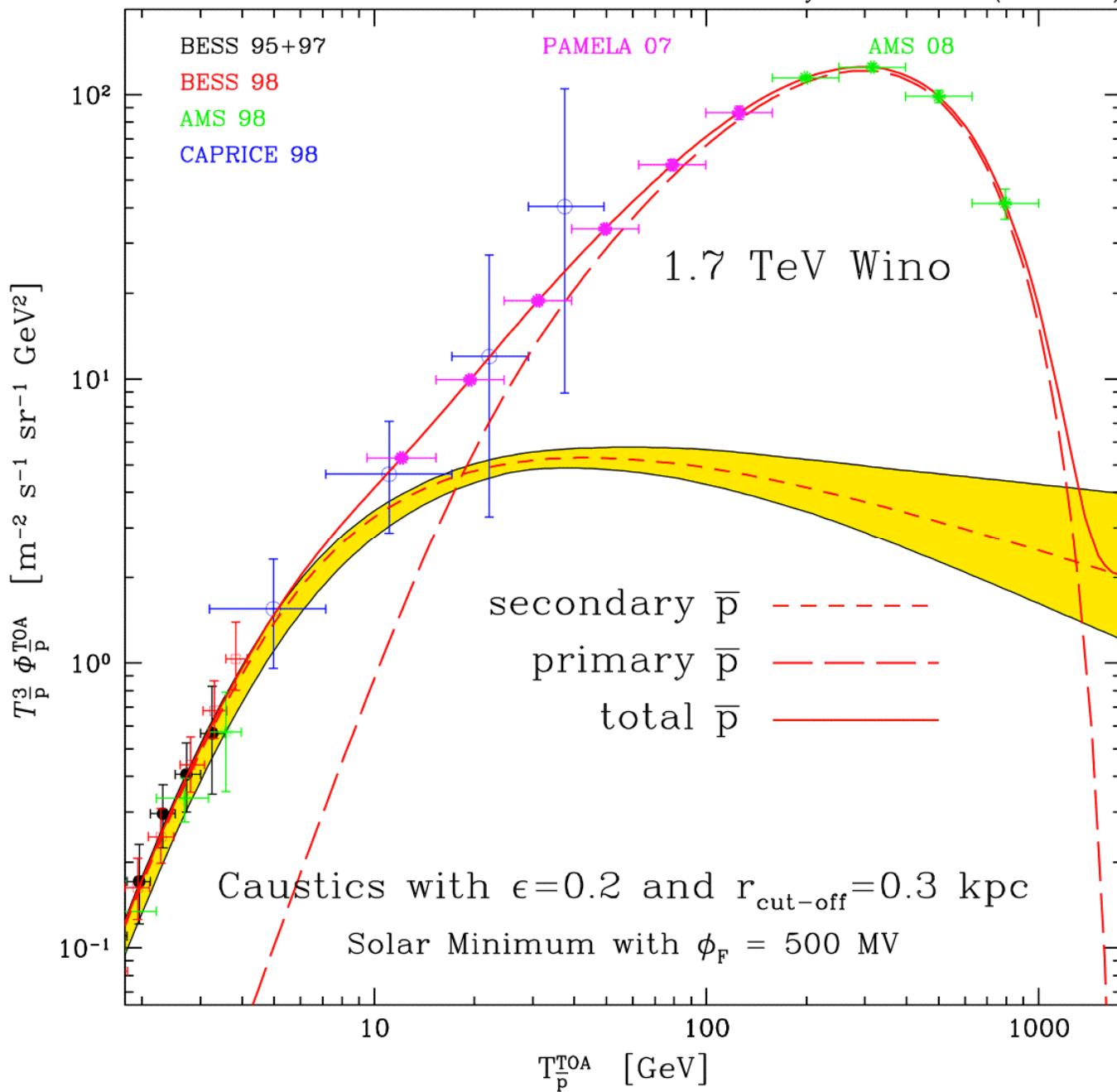
FIG. 7: Expected value and variance of the boost factor of the antiproton signal as a function of kinetic energy, in the case of a dark matter particle with mass $m_\chi = 1 \text{ TeV}$ and for different \bar{p} propagation configurations.

Antiprotons injected at $E_S \simeq E$



Conclusions and perspectives

- Our predictions of the DM indirect signals are plagued by large uncertainties. We need observations ! In particular CREAM for δ and $^{10}\text{Be}/^{9}\text{Be}$ for K_0 .
- The 10 GeV to 1 TeV antiproton window is promising. Will the Caprice excess be confirmed by PAMELA and AMS?
- A few particle physics models lead to a visible signal. In general, large boost factors are needed though.
- The DM mini-spike scenario is associated to very large boost factors with a signal exceeding – on average – the observations. However, beware of the large variance !
- In the case of caustics, the signal is enhanced by a factor of ~ 30 if $r_{\text{cut off}}$ is as small as 0.3 kpc. But is it reasonable ?



Space diffusion dominates in the master equation

$$V_C \partial_z \Psi - K \Delta \Psi + \partial_E \{ b^{\text{loss}}(E) \Psi - K_{EE}(E) \partial_E \Psi \} = Q$$

Poisson equation $K \Delta \Psi + Q = 0$



Long range with $G_{\bar{p}}^{\text{3D}}(r) = \frac{Q}{4\pi Kr}$

- Evaporation at the vertical boundaries $\pm L$
- Leakage at the radial boundaries $R = 20$ kpc
- Evaporation from convective wind V_C
- Annihilations inside the MW gaseous disk
- Energy losses and mild diffusive reacceleration

4) The hidden realm of Vlasovia and DM caustics

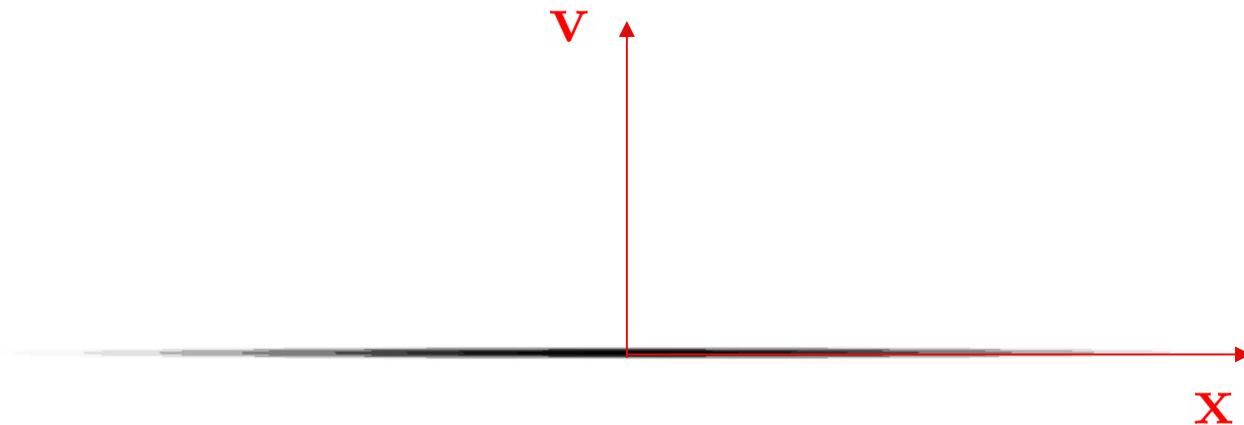
The **Vlasov** equation holds in phase space

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla_{\mathbf{x}} f - \nabla_{\mathbf{x}} \Phi \cdot \nabla_{\mathbf{v}} f = 0 \text{ (no collision)}$$

$$\Delta \Phi = 4\pi G \int d^D \mathbf{v} f(\mathbf{x}, \mathbf{v}, t) \text{ (Poisson equation)}$$

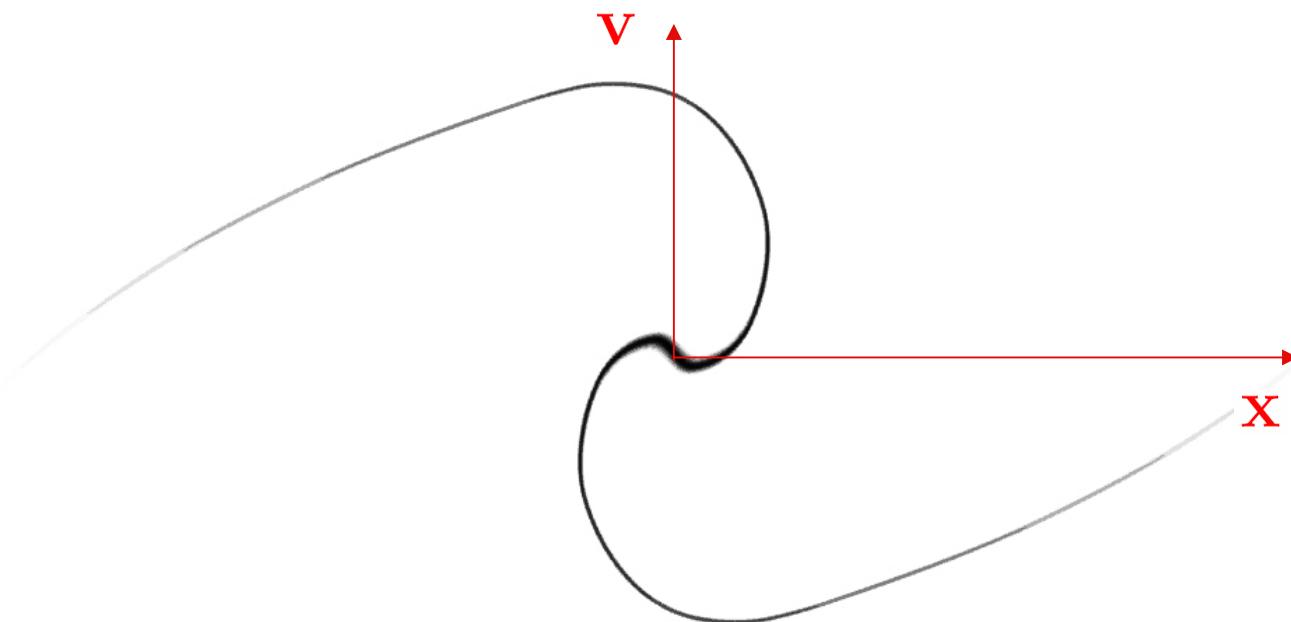
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Cold Initial Conditions

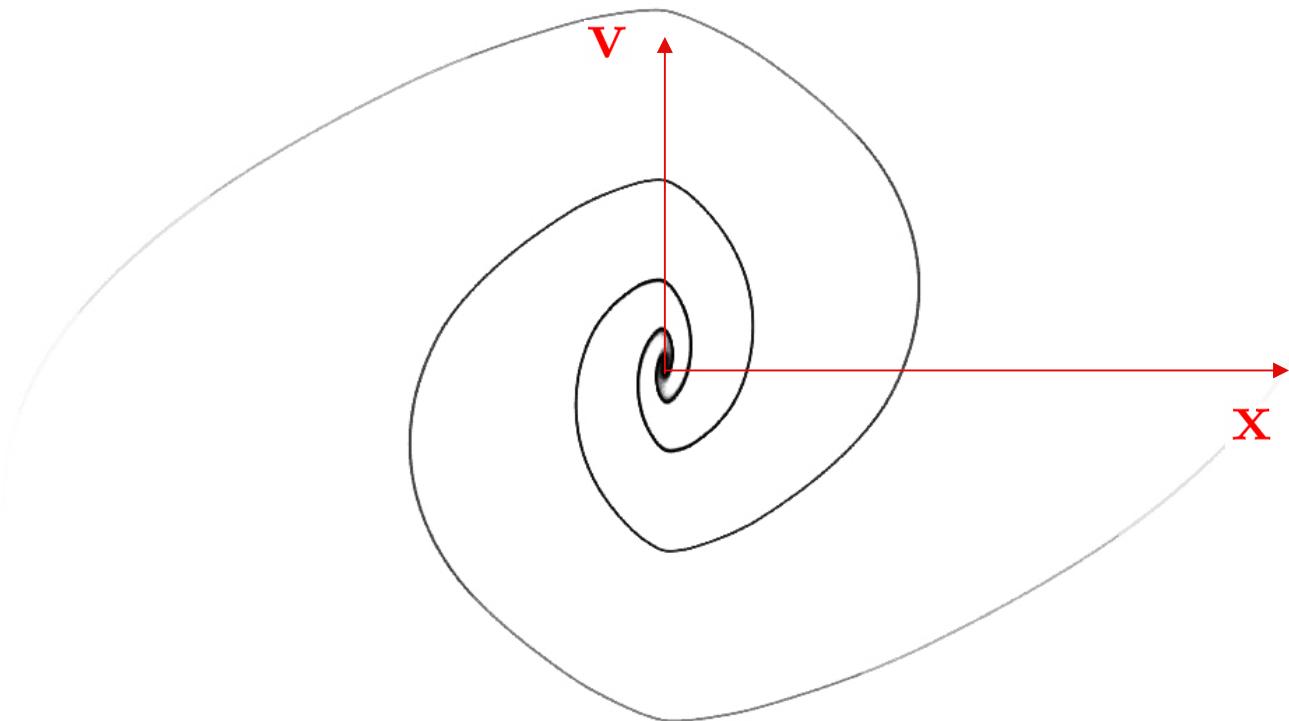


Evolution under 1-D gravity in phase space

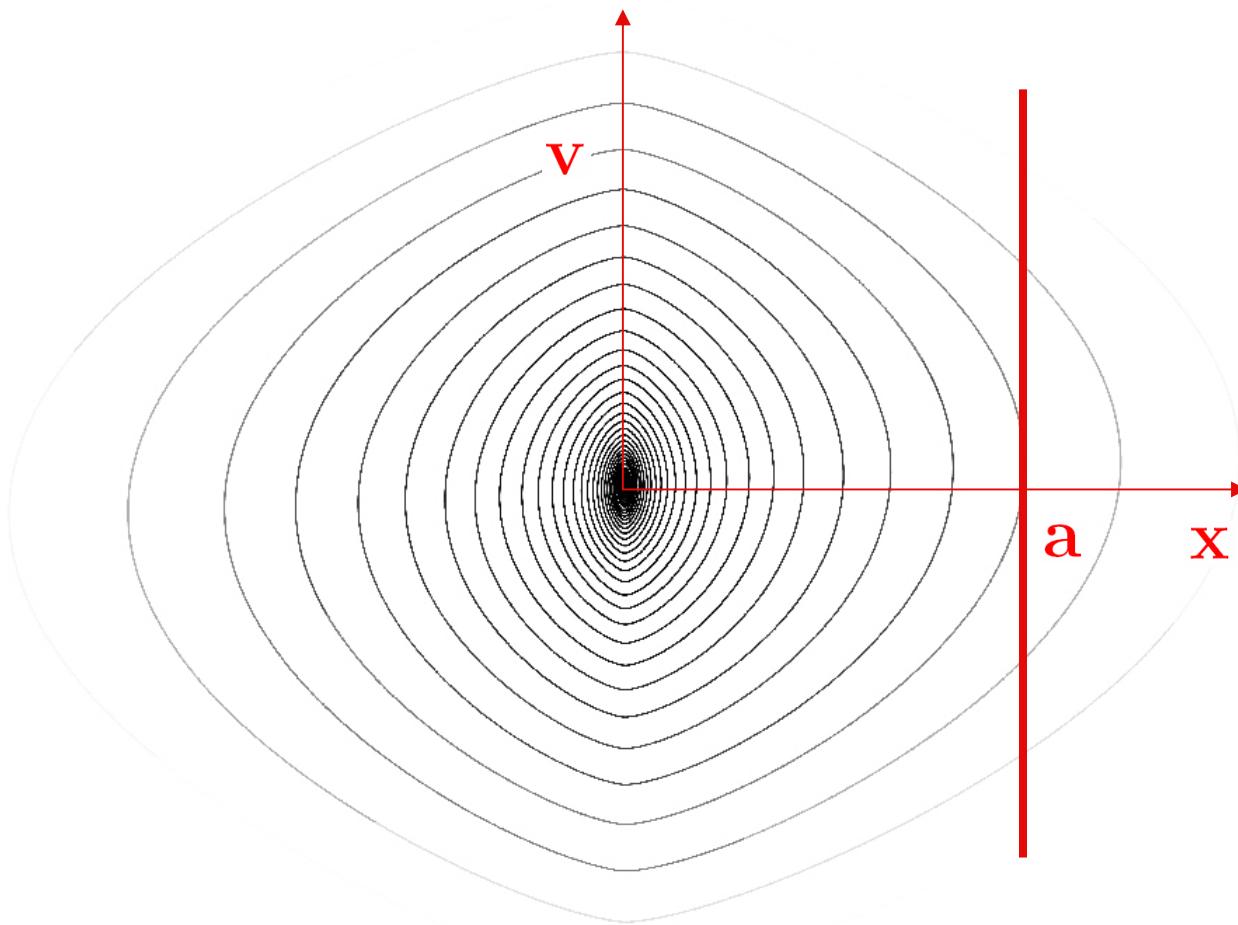
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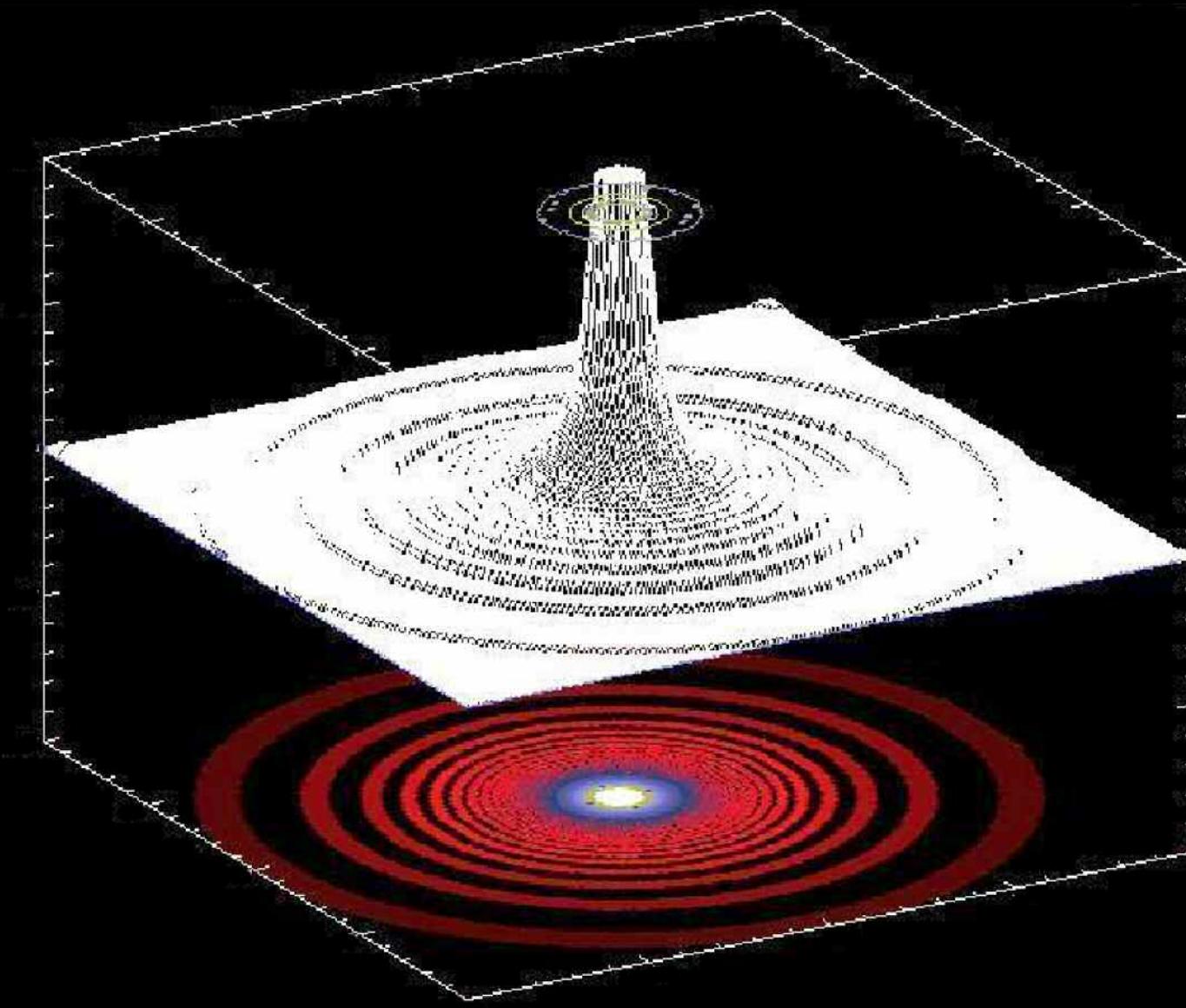


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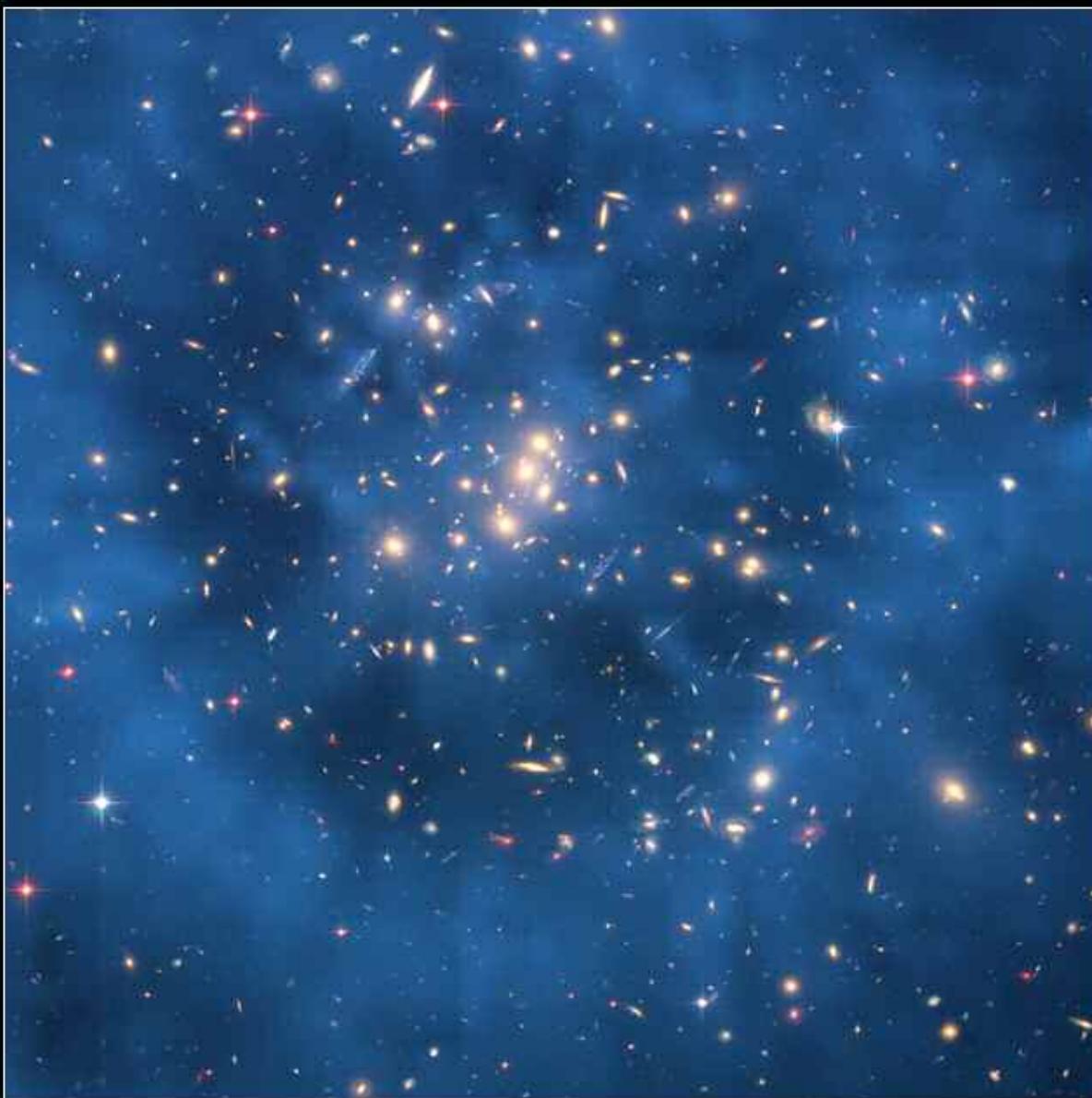


$$\rho(\mathbf{a}) = \int d\mathbf{v} f(\mathbf{a}, \mathbf{v}) \text{ is infinite !}$$

Self-similar infall model – E. Bertschinger, ApJ **58** (1985) 39

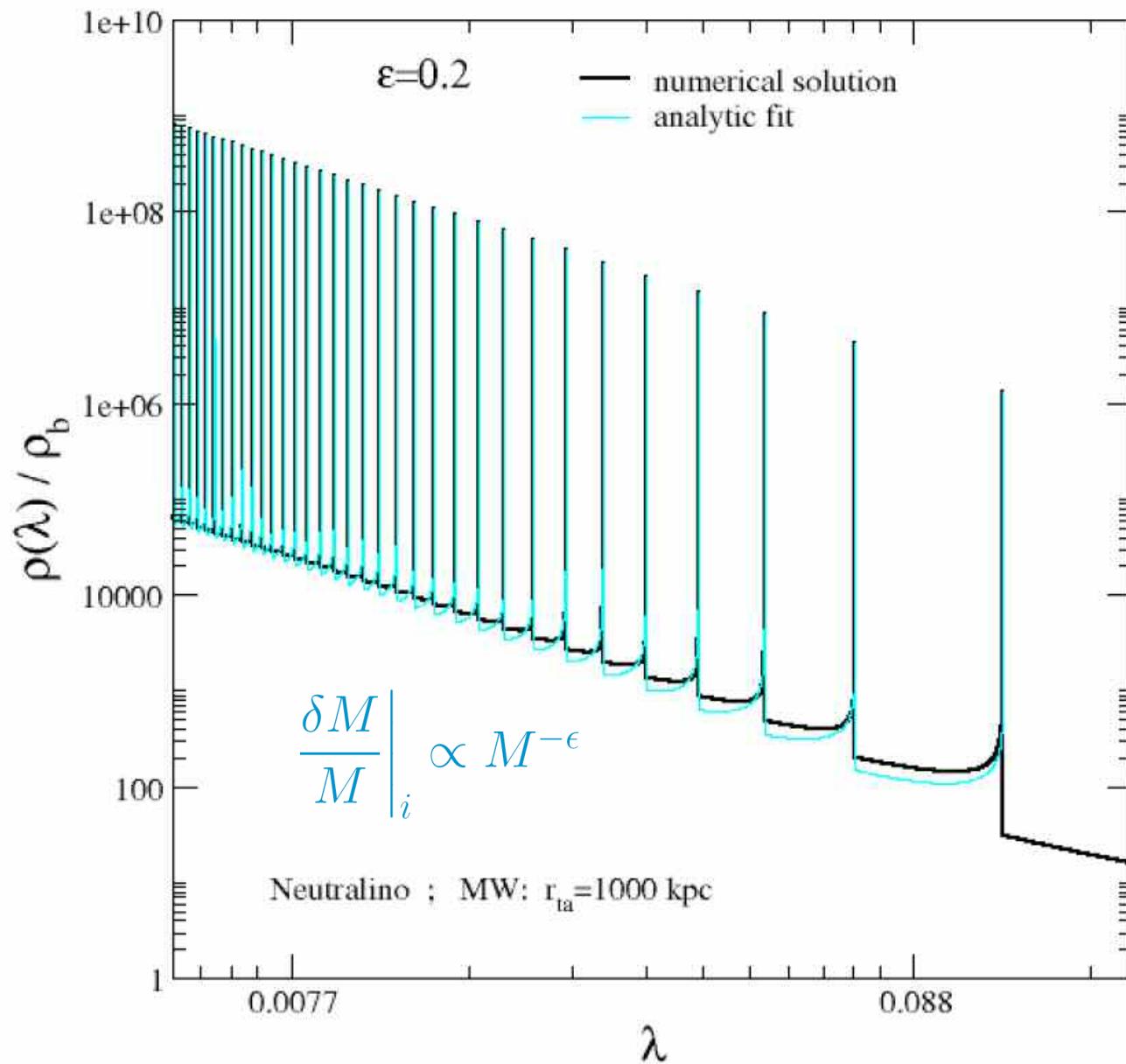


Dark Matter Ring in Cl 0024+17 (ZwCl 0024+1652) HST • ACS/WFC



NASA, ESA, and M.J. Jee (Johns Hopkins University)

STScI-PRC07-17b



λ_D is much larger than the intra-caustic distance

